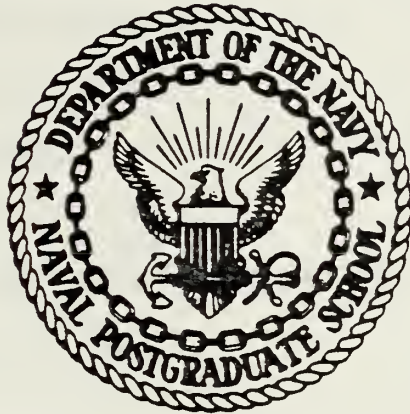


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THESIS

TIME INTEGRATION OF UNSTEADY TRANSONIC FLOW
TO A STEADY STATE SOLUTION BY
THE FINITE ELEMENT METHOD

by

Raymond John Nichols Jr.

March 1977

Thesis Advisor:

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REPORT DOCUMENTATION PAGE		READ INSTRUCTIONS BEFORE COMPLETING FORM
1. REPORT NUMBER	2. GOVT ACCESSION NO.	3. RECIPIENT'S CATALOG NUMBER
4. TITLE (and Subtitle) Time Integration of Unsteady Transonic Flow to a Steady State Solution by the Finite Element Method		5. TYPE OF REPORT & PERIOD COVERED Master's Degree; March 1977
7. AUTHOR(s) Raymond John Nichols Jr.		6. PERFORMING ORG. REPORT NUMBER
9. PERFORMING ORGANIZATION NAME AND ADDRESS Naval Postgraduate School Monterey, California 93940		8. CONTRACT OR GRANT NUMBER(s)
11. CONTROLLING OFFICE NAME AND ADDRESS Naval Postgraduate School Monterey, California 93940		10. PROGRAM ELEMENT, PROJECT, TASK AREA & WORK UNIT NUMBERS
14. MONITORING AGENCY NAME & ADDRESS (if different from Controlling Office) Naval Postgraduate School Monterey, California 93940		12. REPORT DATE March 1977
		13. NUMBER OF PAGES 110
		15. SECURITY CLASS. (of this report) Unclassified
		15a. DECLASSIFICATION/DOWNGRADING SCHEDULE
16. DISTRIBUTION STATEMENT (of this Report) Approved for public release; distribution unlimited.		
17. DISTRIBUTION STATEMENT (of the abstract entered in Block 20, if different from Report)		
18. SUPPLEMENTARY NOTES		
19. KEY WORDS (Continue on reverse side if necessary and identify by block number) Unsteady Transonic Small Perturbation Flow Steady Transonic Small Perturbation Flow Finite Element Method Method of Weighted Residuals Converging-Diverging Nozzle		
20. ABSTRACT (Continue on reverse side if necessary and identify by block number) A finite element method was applied to the unsteady transonic small disturbance equation and integrated until the solution converged to the steady state for a thin non-lifting airfoil. The method of weighted residuals was used to formulate the finite element equations and Houbolt's method of central differencing in time was used to integrate these equations.		

Unclassified

SECURITY CLASSIFICATION OF THIS PAGE(When Data Entered)

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Unclassified

SECURITY CLASSIFICATION OF THIS PAGE(When Data Entered)

Time Integration of Unsteady Transonic Flow
to a Steady State Solution by
the Finite Element Method

by

Raymond John Nichols Jr.
Lieutenant, United States Navy
B.S., Pennsylvania State University, 1970

Submitted in partial fulfillment of the
requirements for the degree of

MASTER OF SCIENCE IN AERONAUTICAL ENGINEERING

from the

NAVAL POSTGRADUATE SCHOOL

March 1977

ABSTRACT

A finite element method was applied to the unsteady transonic small disturbance equation and integrated until the solution converged to the steady state for a thin non-lifting airfoil. The method of weighted residuals was used to formulate the finite element equations, and Houbolt's method of central differencing in time was used to integrate these equations.

A secondary investigation applied the steady transonic small disturbance equations to a converging-diverging nozzle.

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I. INTRODUCTION

Transonic inviscid flows past a smooth airfoil may be expressed in terms of the velocity potential ϕ satisfying the transonic small disturbance equation,

$$(1 - M_\infty^2 - M_\infty^2(\gamma + 1)\phi_x)\phi_{xx} + \phi_{yy} = 0 \quad (1)$$

This equation presents two major difficulties, 1) it is non-linear and 2) it is of the mixed hyperbolic-elliptic type. Analytical solutions to non-linear equations are difficult to obtain. One must normally resort to numerical methods. When the coefficient $(1 - M_\infty^2 - M_\infty^2(\gamma + 1)\phi_x)$ in equation 1 is negative, the flow is supersonic and the equation is called hyperbolic; otherwise the flow is subsonic and the equation is elliptic. The forms of the two solutions are fundamentally different. Hyperbolic equations admit both discontinuities, which propagate only in characteristic directions, and the presence of shock fronts. Elliptic equations, on the other hand, require that the dependent variables and their derivatives be continuous and that a change in any part of the flow field affects every other part. Many non-linear elliptic equations are solved by appropriate relaxation iteration techniques by casting the equation in Poisson's form with the non-linearity acting as the driving force. Solutions to hyperbolic equations are usually obtained by the method of characteristics or by

finite difference marching techniques which use an artificial viscosity to represent the average jump conditions across the shock wave. In mixed supersonic and subsonic flows, normal numerical procedures break down because the boundary between the two regions is not known a priori.

Finite element numerical techniques have evolved as a powerful tool in obtaining approximate solutions to a wide variety of engineering problems, particularly ones with Neumann-Dirichlet boundary conditions, i.e., elliptical problems. They offer several outstanding advantages. Some of these are:

- 1) Non-homogeneous problems may be treated with relative simplicity.
- 2) Complex geometries may be modeled with relative ease since the elements can be graded in size and shape to follow boundaries of arbitrary shape.
- 3) Once the finite element model has been established, a variety of problems can be solved by supplying the computer with the appropriate data.

Chan and Brashears [Ref. 5] developed a finite element computer program for steady transonic flow over a non-lifting airfoil. This program uses the least squares method of weighted residuals to approximate equation 1 by a system of algebraic equations, and assembles the equations in a special way to account for the hyperbolic region of flow. This technique prevents the influence of downwind nodes from propagating upstream in the supersonic region.

The purpose of this thesis is to investigate the possibility of speeding the convergence of a solution by transforming the steady transonic equation to the unsteady equation

and integrating over time until the time dependent terms vanish and to extend the program of Ref. [5] to the transonic region of a converging-diverging nozzle.

II. DISCUSSION OF THE FINITE ELEMENT APPROACH

In a continuum problem of any dimension, the field variable, whether it is velocity potential, velocity, temperature, displacement or some other quantity, possesses infinitely many values because it is a function of each generic point in the solution region. Consequently, the problem is one of an infinite number of unknowns. The finite element approach subdivides the solution domain into a finite number of subdomains called elements and expresses the unknown field variable in terms of assumed approximating functions within each element. The approximating functions are sometimes called interpolating functions and are defined in terms of the values of the field variables at specified points called nodes or nodal points. Nodes usually lie on the element boundaries where adjacent elements are considered to be connected. In addition to boundary nodes, an element may also contain interior nodes. The nodal values of the field variable and the interpolating function for the elements completely define the behavior of the field variable within the elements. Once these unknowns are found, the interpolating functions define the field variable throughout the assemblage of the element.

Clearly, the nature of the solution and the degree of accuracy of the approximation depends not only on the number and size of the elements used but also on the interpolating

functions selected. Interpolating functions may not be chosen arbitrarily because certain compatibility conditions must be satisfied. Often such functions are selected so that the field variable or its derivatives are continuous across adjoining element boundaries. Once the problem is formulated in terms of individual elements, the contributions of each element may be assembled to define the entire solution domain. This means, for example, that if we are treating a problem in stress analysis we can find the force-displacement or stiffness characteristics of each element and then assemble the elements to determine the stiffness of the whole structure. Finite element solutions are not, of course, restricted to structures problems, but the matrix of equations defined by the interpolating functions and the nodal field variables is still referred to as the stiffness matrix regardless of the field variable in the problem.

Solutions to continuous problems by the finite element approach always follow a systematic step-by-step process. This process is completely general to the finite element method and it is outlined below. [Ref. 4]

1. Discretize the continuum.

The first step is to divide the solution domain into elements. A variety of element shapes may be used and one or more different element shapes may be used in the same region. The type and number of the elements used in a given problem are a matter of engineering judgement.

2. Select the interpolating functions.

The next step is to choose the type of interpolating function to represent the variation of the field variable over the element. The field variable may be a scalar, a vector, or a higher order tensor. Often polynomials are selected as interpolating functions for they are easy to integrate and differentiate. The degree of polynomial chosen depends on the number of nodes and the nature and number of the unknowns and the continuity requirements imposed at the nodes and the element boundaries. The unknown quantities at the nodes may be assigned to the field variable and their derivatives.

3. Find the element properties.

After the elements and their interpolating functions have been selected, the matrix of algebraic equations which express the properties of the individual elements must be determined. Several methods are available for this task. These are: the direct approach, the variational approach, the method of weighted residuals and the energy balance approach. Reference [4] is a good source of information on the various techniques.

4. Assemble the element properties to obtain the system equations.

To find the properties of the over-all system, the matrix equations expressing the behavior of each element must be

added to the matrix equation of all other elements. The basis for this assembly procedure stems from the fact that connecting elements have common nodes and the field variable must be the same for each element sharing that node.

At this point the boundary conditions for the system of equations must be accounted for and the system of equations must be modified before it is ready for solution.

5. Solve the system of equations.

The assembly process of step 4 produces a set of simultaneous equations which can be solved to obtain the unknown field variables. Linear equations have a number of standard solution techniques readily available, but solutions to non-linear equations are more difficult to obtain.

6. Make additional computations if desired.

Important parameters, such as pressure coefficient in aerodynamic problems, may now be calculated from the values of the field variables.

III. THEORY AND BASIC EQUATIONS

A. STEADY TRANSONIC FLOW

Chan et al. [Ref. 5] developed an algorithm to analyze steady transonic flow over non-lifting thin airfoils. Boundary layer effects were ignored and the imbedded shock wave was assumed to be weak. These assumptions are consistent with transonic small disturbance theory which can be expressed mathematically by the following expressions.

$$(1 - M_{\infty}^2 - M_{\infty}^2(\gamma + 1)\phi_x)\phi_{xx} + \phi_{yy} = 0 \quad (1)$$

Boundary conditions -

$$\nabla \cdot \phi = 0 \quad \text{at infinity} \quad (2)$$

$$v = (1 + u)dg/dx \quad \text{on the airfoil} \quad (3)$$

$$u = 0 \quad \text{on the line of symmetry} \quad (4)$$

where $g(x,y)$ defines the airfoil and dg/dx describes the airfoil slope.

The above expressions appear in their dimensionless form where ϕ = perturbed velocity potential and the perturbed velocity components in the x and y directions are respectively defined as

$$u = \phi_x$$

$$v = \phi_y$$

M_∞ = freestream Mach number and γ = ratio of specific heats which for air is taken to be 1.4. The physical coordinates x' and y' and the velocity potential ϕ' are related to the dimensionless quantities by

$$x = x'/c, \quad y = y'/c, \quad \phi = \phi'/cU_\infty$$

where c is the chordlength of the airfoil and U_∞ is the free-stream velocity.

Once the flowfield solution is determined in terms of the perturbed velocities, the secondary unknowns are subsequently calculated. These include:

$$a = \left[\frac{\gamma-1}{2} (U_\infty^2 - V^2) + \left(\frac{U_\infty}{M} \right)^2 \right]^{\frac{1}{2}} \quad (5)$$

$$M = \frac{V}{a} \quad (6)$$

$$\frac{P}{P_0} = \frac{1}{\left[1 + \frac{\gamma-1}{2} M^2 \right]^{\gamma/(\gamma+1)}} \quad (7)$$

$$C_p = - \left[\frac{2u}{U_\infty} + (1-M_\infty^2) \frac{u^2}{U_\infty^2} + \frac{v^2}{U_\infty^2} \right] \approx - \frac{2u}{U_\infty} \quad (8)$$

In the above, $U_\infty = 1$ is the normalized freestream velocity, a = local sound speed, p = local static pressure, M = local Mach number, V = total velocity, P_0 = stagnation pressure and C_p = pressure coefficient.

B. UNSTEADY TRANSONIC FLOW

Unsteady transonic inviscid flow may be expressed in terms of the velocity potential $\phi(x,y,t)$ to a first order approximation by

$$(1 - M_\infty^2 - M_\infty^2(\gamma + 1)\phi_x)\phi_{xx} + \phi_{yy} - 2M_\infty^2\phi_t - M_\infty^2\phi_{tt} = 0 \quad (9)$$

which has the same non-linear coefficient retained for steady transonic flow in Equation 1.

Boundary conditions require that the disturbances vanish far from the airfoil,

$$\phi_x = 0 \quad \phi_y = 0 \quad \text{at infinity} \quad (10)$$

and that the flow remain attached to the body. Let $B(x,y,t) = 0$ define the body at any instant. The surface tangency restraint may now be expressed by the substantial derivative DB/DT vanishing.

$$DB/DT = B_t + (1 + \phi_x)B_x + \phi_y B_y \quad (11)$$

If the body remains stationary $B_t = 0$, and the tangency condition becomes the same as in steady flow, namely

$$v = (1 + u)dg/dx \quad (12)$$

where dg/dx represents the airfoil slope.

The pressure coefficient for isentropic unsteady compressible flow is defined by

$$C_p = \frac{2}{\gamma M_\infty^2} \left\{ \left[1 - \frac{\gamma-1}{2} M_\infty^2 (2\phi_x + 2\phi_t + \phi_x^2 + \phi_y^2) \right]^{\gamma/(\gamma-1)} - 1 \right\} \quad (13)$$

Expanding by the binomial expansion and retaining only the first order terms gives,

$$C_p = -2\phi_x - 2\phi_t \quad (14)$$

IV. FINITE ELEMENT FORMULATION

The method of weighted residuals is a technique for approximating solutions to linear or non-linear partial differential equations and it is the basis for the finite element formulation of the transonic small disturbance equation (Equation 1). This procedure involves assuming the general functional behavior of the field variable which would approximately satisfy the basic equation and boundary conditions. Substituting this approximation into the original differential equation results in some error called a residual. A system of algebraic equations results when a weighted average of the residual is forced to vanish as it is averaged over the entire domain.

A. STEADY FLOW

The approximate solution to equation 1 is assumed to be

$$\hat{\phi} = N_i \phi_i \quad (15)$$

where N_i are the interpolating functions which exhibit the behavior of equation 1 and ϕ_i are the undetermined parameters at the nodal points.

When $\hat{\phi}$ is substituted into equation 1, the resulting residual is

$$R = [1 - M_\infty^2 - M_\infty^2(\gamma + 1)N_{k,x}\phi_k]N_{j,xx} + N_{j,yy} \quad (16)$$

The weighted average is determined by multiplying the residual R by m linearly independent weighting functions W_i and integrating over the elemental domain. Forcing this residual to vanish yields,

$$\iint W_i R dA = 0 \quad i = 1 \text{ to } m$$

Chan et al. [Ref. 5] found that when the weighting function W_i for equation 1 is chosen to be $\partial R / \partial \phi_i$ the resulting matrix is positive definite and well conditioned. This choice of weighting functions is referred to as the method of least squares because it is equivalent to minimizing the square of the residuals summed over the domain with respect to the undetermined parameters. That is,

$$\begin{aligned} \chi &= \iint R^2 dA \\ \partial \chi / \partial \phi_i &= \iint \partial R / \partial \phi_i R dA = 0 \end{aligned} \quad (17)$$

Integrating over the domain produces the system of algebraic equations

$$S_{ij} \phi_j = 0 \quad (18)$$

where the elemental matrix S_{ij} is defined as

$$S_{ij} = \iint Q_j P_i dA \quad (19)$$

With Q_j and P_i equal to

$$Q_j = [1 - M_\infty^2 - M_\infty^2(\gamma + 1)N_{k,x}\phi_k] N_{j,xx} + N_{j,yy}$$

$$P_i = Q_i - M_\infty^2(1 + \gamma)N_{k,xx}\phi_k N_{i,x}$$

B. UNSTEADY FLOW

Development of the unsteady flow finite element equations is similar to the procedure used to formulate the finite element equations for steady transonic flow. The least squares method of weighted residuals is again used but ϕ is now a function of time as well as the spatial coordinates x and y .

The approximate solution has the form,

$$\hat{\phi} = N_i(x,y)\phi_i(t) \quad (20)$$

Substituting $\hat{\phi}$ in the unsteady transonic small disturbance equation, the weighted residual becomes,

$$\chi = \iint (R_1 + R_2 + R_3)^2 dA \quad (21)$$

where

$$R_1 = \{ [1 - M_\infty^2 - M_\infty^2(\gamma + 1)\phi_k N_{k,x}] N_{j,xx} + N_{j,yy} \} \phi_j$$

$$R_2 = -2M_\infty^2 N_{j,x} \dot{\phi}_j$$

$$R_3 = -M_\infty^2 N_j \ddot{\phi}_j$$

Expanding equation 21 yields

$$\chi = \iint [R_1^2 + R_2^2 + R_3^2 + 2R_1R_2 + 2R_2R_3] dA \quad (22)$$

and minimizing with respect to the undetermined parameters ϕ_i the following system of algebraic equations result,

$$\phi_j = 0 = 2 \iint \partial R_i / \partial \phi_j [R_1 + R_2 + R_3] dA$$

$$\partial R_i / \partial \phi_j = P_i$$

where P_i has been previously defined in the steady finite element formulation

The above equation may be rewritten in the form

$$S_{ij}\phi_j + SC_{ij}\dot{\phi}_j + SM_{ij}\ddot{\phi}_j = 0 \quad (23)$$

The stiffness matrix S_{ij} is the same as that developed for the steady transonic equation and the damping (SC_{ij}) and mass (SM_{ij}) matrices are defined below.

$$SC_{ij} = -\iint M_\infty^2 N_{j,x} P_i dA \quad (24)$$

$$SM_{ij} = -\iint M_\infty^2 N_j P_i dA \quad (25)$$

V. ELEMENT DESCRIPTION AND ASSEMBLY OF EQUATIONS

A. ELEMENT DESCRIPTION

The basic element used in the finite element program is the non-conforming cubic triangular element developed by Bazeley et al. [Ref. 2]. Also used in the program are the quadrilateral and trapezoidal elements constructed from these triangular elements. These three types of elements can be mixed and used freely in the entire flow region except that only trapezoids should be used to cover the supersonic and mixed region in order to enact the special assembly procedures required by the hyperbolic equation which describes the flow in that region.

The basic triangular element is shown in Fig. 1, which at each vertex has the velocity potential and the velocity components as the undetermined parameters. This type of element was chosen because both Dirichlet and Neumann boundary conditions can be imposed with equal convenience. In addition, because velocity components are used as primary unknowns secondary parameters, such as Mach number and pressure coefficient can be calculated directly without resorting to numerical differentiation, which would produce additional errors.

In the element, the approximate solution is assumed as

$$\hat{\phi} = N_i \phi_i \quad (i = 1 \text{ to } 9)$$

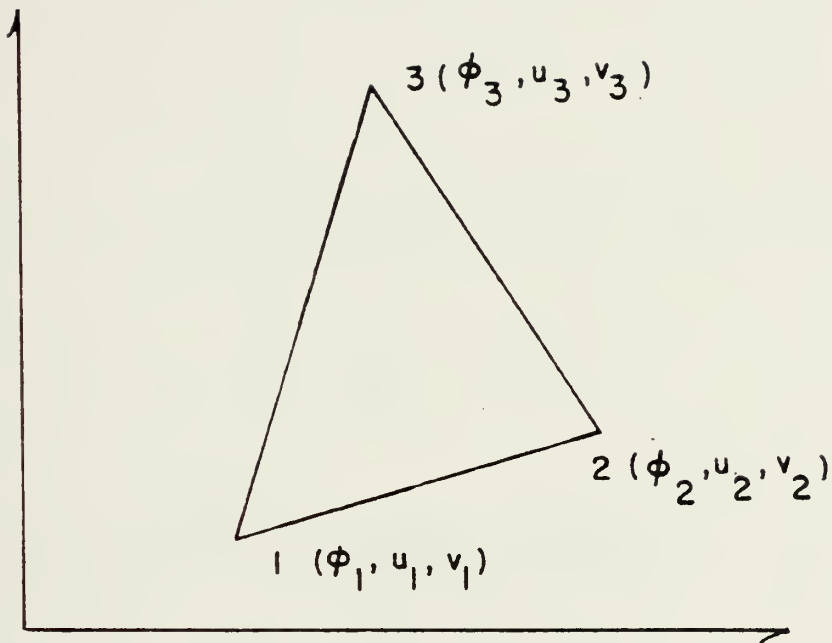


Figure 1 - Triangular Element

In which ϕ_i 's are the nine undetermined parameters of ϕ and N_i are the interpolation functions which are expressed in terms of the area coordinates.

The shape or interpolating functions and their first and second derivatives are defined below.

Letting,

$$a_i = x_j y_k - x_k y_j$$

$$b_i = y_j - y_k$$

$$c_i = x_k - x_j$$

$$\Delta = \text{area of triangle 1-2-3} = (b_j c_k - b_k c_j)/2$$

$$\alpha = 0.5 (c_k - c_j)$$

$$\beta = 0.5 (b_j - b_k)$$

$$H = \zeta_k \zeta_j \zeta_i$$

$$H_x = b_i \zeta_j \zeta_k + b_j \zeta_k \zeta_i + b_k \zeta_i \zeta_j$$

$$H_y = c_i \zeta_j \zeta_k + c_j \zeta_k \zeta_i + c_k \zeta_i \zeta_j$$

$$H_{xx} = 2(\zeta_k b_j b_k + \zeta_j b_k b_k + \zeta_k b_i b_j)$$

$$H_{yy} = 2(\zeta_i c_j c_k + \zeta_j c_k c_i + \zeta_k c_i c_j)$$

with $i = (1,2,3), k = (3,1,2)$ then one has for $l = (1,4,7),$
 $i = (1,2,3)$

$$N_1 = \zeta_i^2 (3 - 2\zeta_i) + 2H$$

$$N_{1,x} = [6b_i \zeta_i (1 - \zeta_i) + 2H_x]/2\Delta$$

$$N_{1,y} = [6c_i \zeta_i (1 - \zeta_i) + 2H_y]/2\Delta$$

$$N_{1,xx} = [6b_i^2 (1 - 2\zeta_i) + 2H_{xx}] / (2\Delta)^2$$

$$N_{1,yy} = [6c_i^2 (1 - 2\zeta_i) + 2H_{yy}] / (2\Delta)^2$$

for $1 = (2, 5, 8)$, $1 = (1, 2, 3)$

$$N_1 = \zeta_i^2 (c_k \zeta_j - c_j \zeta_k) + \alpha H$$

$$N_{1,x} = [2b_i \zeta_i (c_k \zeta_j - c_j \zeta_k) + 2\zeta_i^2 + \alpha H_x] / 2\Delta$$

$$N_{1,y} = [2c_i \zeta_i (c_k \zeta_j - c_j \zeta_k) + \alpha H_y] / 2\Delta$$

$$N_{1,xx} = [2b_i^2 (c_k \zeta_j - c_j \zeta_k) + 4b_i (2\Delta) \zeta_i + \alpha H_{xx}] / (2\Delta)^2$$

$$N_{1,yy} = [2c_i^2 (c_k \zeta_j - c_j \zeta_k) + \alpha H_{yy}] / (2\Delta)^2$$

for $1 = (3, 6, 9)$, $1 = (1, 2, 3)$

$$N_1 = \zeta_i^2 (b_j \zeta_k - b_k \zeta_j) + \beta H$$

$$N_{1,x} = [2b_i \zeta_i (b_j \zeta_k - b_k \zeta_j) + \beta H_x] / 2\Delta$$

$$N_{1,y} = [2c_i \zeta_i (b_j \zeta_k - b_k \zeta_j) + 2\zeta_i^2 + \beta H_y] / 2\Delta$$

$$N_{1,xx} = [2b_i^2 (b_j \zeta_k - b_k \zeta_j) + \beta H_{xx}] / (2\Delta)^2$$

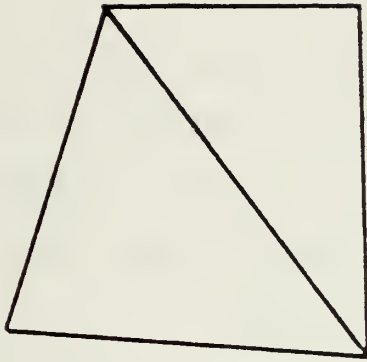
$$N_{1,yy} = [2c_i^2 (b_j \zeta_k - b_k \zeta_j) + 4c_i (2\Delta) \zeta_i + \beta H_{yy}] / (2\Delta)^2.$$

Quadrilateral and trapezoidal elements as shown in Fig. 2 are also used in the present program, the former is used in the subsonic region and the latter in the mixed and supersonic region. The element matrix for the quadrilateral element is obtained by combining appropriately the matrices for two triangles, while the matrix for trapezoidal element is obtained by averaging contributions from two left-running and two right-running triangles. The averaging process is designed to remove the bias effects inherent in the quadrilaterals used.

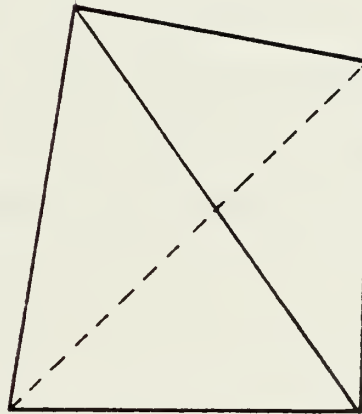
B. ASSEMBLY OF EQUATIONS

Straightforward application of the finite element assembly technique to transonic flows would fail (the solution diverges) because this would allow disturbances to propagate upwind in the supersonic region of flow where the governing equation is hyperbolic. Hyperbolic equations have a time-like dependency in that the solution at the downwind station is affected by the upwind station but not vice-versa. Assembly techniques for a transonic flow finite element program must take into account this time-like dependency. If the x-axis is taken as a time-like direction in the supersonic region, the element matrix may be assembled in a way similar to a backward finite difference operator, which has been successful in solving hyperbolic equations.

Consider the rectangular element sketched below with the upwind station I and the downwind station II, each having two nodal points with the element type chosen. The element matrices can be constructed in the usual manner.



a. Quadrilateral
Element



b. Trapezoidal
Element

Figure 2 - Quadrilateral and Trapezoidal Elements

However, before assembling the element matrix into the system matrix the non-linear coefficient of equation 1 is evaluated.

$$C = 1 - M_{\infty}^2 - M_{\infty}^2 (\gamma + 1) u$$

The sign of the coefficient being positive, zero, or negative defines the equation as elliptic, parabolic, or hyperbolic. If C is non-positive for all nodes in the element, the rows representing the improper downwind influence on the solution at an upwind station are ignored during assembly. This feature is automatically applied in the program requiring only a little care in arrangement of the nodes of the element. In the anticipated supersonic region, element node points should be arranged in the order as indicated in figure 3, starting with the upper left corner and proceeding in the counter-clockwise direction. In the elliptic region, i.e., where the coefficient is positive, no special assembly technique is invoked.

C. ITERATIVE PROCEDURES

With the equations assembled and the proper boundary conditions imposed, the system of non-linear algebraic equations is solved by iterative procedures in the form

$$S_{ij}(\tilde{\phi})\phi_j^{(n)} = 1_i \quad (23)$$

to solve for the solution in the n^{th} iteration. The function $\tilde{\phi}$ is defined as

$$\tilde{\phi} = \theta\phi^{(n-1)} + (1-\theta)\phi^{(n-1)}$$

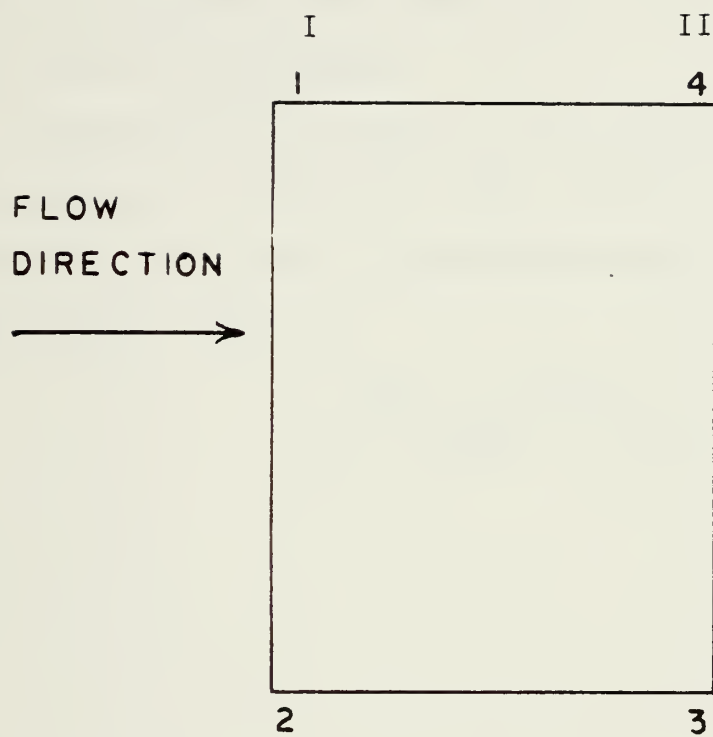


Figure 3 - Nodal Arrangement for Supersonic Region

in which the under-relaxation factor θ is in the range $0 < \theta \leq 1$. For subsonic flow $\theta = 1$, which is simply a successive approximation, yields good results, but it is necessary to under-relax somewhat with θ approximately .5 for supercritical flow. Generally, a smaller relaxation factor will make the solution more stable but it will tend to slow down the rate of convergence.

Equation 23 is subject to the convergence criterion that the change in local Mach number between two successive iterations is less than a prescribed value ϵ at all nodes in the flow field. That is,

$$\left| \frac{M^{(n)} - M^{(n-1)}}{M^{(n)}} \right| \leq \epsilon .$$

VI. INTEGRATION OF UNSTEADY FINITE ELEMENT EQUATIONS

The unsteady transonic small disturbance equation (equation 9) when suitably reduced to a finite element approximation appears in the form,

$$S_{ij}\phi_j + SC_{ij}\dot{\phi}_j + SM_{ij}\ddot{\phi}_j = R_j \quad (26)$$

This equation is analogous to a damped spring mass system, hence S_{ij} , SC_{ij} , and SM_{ij} are respectively referred to as the stiffness, damping, and mass matrices.

Mathematically, equation 26 represents a system of second order differential equations with constant coefficients, which can be solved by standard numerical procedures for differential equations, such as Runge-Kutta or Milne methods. However, this would be a very costly technique if the coefficient matrices are very large. In practical finite element analysis there are a few effective methods which take advantage of the banded matrices usually encountered in finite elements. One such method is the direct numerical integration method.

Direct integration involves a numerical step-by-step procedure aimed at satisfying equation 26 only at discrete time intervals Δt apart and not over all time t . Conceptually, direct integration is a finite element method in space and a finite difference method in time. Examples of direct integration are the central difference method, Houbolt integration and the Wilson method. The first two schemes are finite

difference schemes whereas the latter is a linear acceleration method. Linear acceleration integration assumes a linear variation of acceleration from time t to time $t + \Delta t$.

Central differencing can be very effective in the solution of many dynamic problems especially those that involve a large system of equations. However, this method is unstable for all time steps larger than a critical time step.

Houbolt integration is an implicit finite differencing method related to central differencing, only it has the advantage of being stable for all TIME STEPS.

The Houbolt method was used to integrate the unsteady finite element equations because of this stability.

Houbolt integration uses the following finite difference expansions:

$$\phi_{i,t+\Delta t} = [2\phi_{i,t+\Delta t} - 5\phi_{i,t} + 4\phi_{i,t-\Delta t} - \phi_{i,t-2\Delta t}]/\Delta t^2$$

$$\phi_{i,t+\Delta t} = [11\phi_{i,t+\Delta t} - 18\phi_{i,t} + 9\phi_{i,t-\Delta t} - 2\phi_{i,t-2\Delta t}]/6\Delta t$$

which are two backward-difference formulas with errors of order $(\Delta t)^2$.

The solution of $\phi_{i,t+\Delta t}$ must satisfy equation 26 and at time $t+\Delta t$ equation 26 becomes

$$S_{ij}\phi_{j,t+\Delta t} + SC_{ij}\dot{\phi}_{j,t+\Delta t} + SM_{ij}\ddot{\phi}_{j,t+\Delta t} = 0$$

Substituting the finite difference formulas for $\phi_{j,t+\Delta t}$ and rearranging all known vectors on the right hand side, the solution for $\phi_{j,t+\Delta t}$ is obtained, namely:

$$(a_0^{SM_{ij}} + a_1^{SC_{ij}} + S_{ij})\phi_{j,t+\Delta t} = R_{j,t+\Delta t} \quad (27)$$

$$+ (a_2^{SM_{ij}} + a_3^{SC_{ij}})\phi_{j,t} + (a_4^{SM_{ij}} + a_5^{SC_{ij}})\phi_{j,t-2\Delta t}$$

Where the constant integration coefficients are:

$$a_0 = 2/\Delta t^2$$

$$a_1 = 11/6\Delta t$$

$$a_2 = 5/\Delta t$$

$$a_3 = 3/\Delta t$$

$$a_4 = 2a_0$$

$$a_5 = -a_3/2$$

$$a_6 = a_0/2$$

$$a_7 = a_3/9$$

Equation 27 may be written as

$$SE_{ij}\phi_{j,t+\Delta t} = RE_j \quad (28)$$

where the effective stiffness matrix SE_{ij} and the effective load vector RE_j are defined as:

$$SE_{ij} = S_{ij} + a_0^{SM_{ij}} + a_1^{SC_{ij}}$$

$$RE_j = SM_{ij}(a_2\phi_{j,t} + a_4\phi_{j,t-\Delta t} + a_6\phi_{j,t-2\Delta t}) \\ + SC_{ij}(a_3\phi_{j,t} + a_5\phi_{j,t-\Delta t} + a_7\phi_{j,t-2\Delta t})$$

Accurate knowledge of the vectors $\phi_{j,t-\Delta t}$ and $\phi_{j,t-2\Delta t}$ are required to yield an accurate solution for $\phi_{j,t+\Delta t}$ and

normally the Houbolt integration scheme requires a special starting procedure to determine the initial two vectors $\phi_{j,\Delta t}$ and $\phi_{j,2\Delta t}$. However, since the primary interest of this problem is to integrate the equations until they converge to a steady state solution, it is not necessary to obtain an accurate time history of the flow. Errors induced by the inaccurate starting vectors will vanish as time approaches infinity. Therefore, the starting vectors may be chosen somewhat arbitrarily.

VII. CONVERGING-DIVERGING NOZZLE

S. F. Shen [Ref. 10] demonstrated the feasibility of calculating compressible flows through a converging-diverging (Laval) nozzle by dividing the region of calculations into three patches, a subsonic region, a supersonic region and a transonic one, of course bounded by the other two regions. The locations of the boundaries for each region were chosen arbitrarily provided the sonic line is bracketed by the subsonic and supersonic boundaries.

Two different finite element formulations were used for the subsonic and supersonic regions, but Shen [10] resorted to analytical approximations to cover the transonic patch. This restricted the calculations to nozzles with small throat curvatures because no analytic solutions exist for nozzles with large throat curvatures. It is conceivable that STRANL-II could be adapted to provide a continuous solution throughout all three regions.

Outside the transonic region of flow the governing small disturbance equation is

$$(1 - M_{\infty}^2)\phi_{xx} + \phi_{yy} = 0 \quad (29)$$

This holds for both subsonic and supersonic flow. Comparing equation 29 with equation 1, the transonic small disturbance equation, we notice that only the non-linear coefficient $M_{\infty}^2(\gamma + 1)u$ distinguishes the two equations from each other.

This coefficient becomes negligible when the Mach number becomes less than .8 or greater than 1.2. With this consideration in mind, it was assumed that equation 1 would adequately describe the flow through the nozzle and that the finite element formulation developed for the non-lifting airfoil would apply to the Laval nozzle.

A. BOUNDARY CONDITIONS

Two solutions are possible for a converging-diverging nozzle: 1) Symmetric flow, where the flow is subsonic through the domain, except for a small supersonic region near the wall in the throat, and 2) Asymmetric solution, where the flow accelerates to sonic velocity in the throat and then continues to accelerate to supersonic velocity in the diverging section. Different boundary conditions apply for the two solutions. For the symmetric case, both inlet and exit velocities must be specified. Inlet and exit velocities are equivalent in the subsonic solution. The supersonic solution requires that only the inlet velocities be specified. If the exit velocities are also applied, the problem is overspecified and the solution may not converge.

Velocities at the inlet and exit are not uniform in the y direction, therefore the disturbances cannot be set to zero as in the case of the non-lifting airfoil. Boundary velocities must be calculated by solving equation 29 analytically.

Equation 29 is a linear equation which can be mapped to Laplace's equation,

$$\nabla^2 \phi = 0,$$

by letting $y' = \sqrt{(1 - M_\infty^2)} y$. Laplace's equation is easily solved for the case of the hyperbolic nozzle by transforming from cartesian coordinates to elliptic coordinates. This transformation simplifies the solution because the stream lines must be hyperbolas to follow the nozzle boundary and therefore follow the hyperbolic coordinate $v = \text{constant}$.

If the elliptic coordinates μ and v are chosen such that the curves $\mu = \text{constant}$ are ellipses and the v curves are hyperbolas, then the velocity potential which satisfies Laplace's equation for a hyperbolic nozzle is simply

$$\phi = A\mu$$

where A is a constant of integration. The stream function is

$$\psi = Av$$

The transformation $w = \mu + iv = \cosh^{-1}(2z/a)$ gives rise to the elliptic coordinates

$$y = 1/2 a \cosh \mu \cos v, \quad x = 1/2 a \sinh \mu \sin v$$

$$r = 1/2 a [\cosh^2 \mu - \sin^2 v]$$

$$r_1 = \sqrt{(y + a/2)^2 + x^2}$$

$$r_2 = \sqrt{(y - a/2)^2 + x^2}$$

Solving for μ and v produces

$$\mu = \cosh^{-1} [(r_1 + r_2)/a]$$

$$v = \cos^{-1} [(r_1 - r_2)/a]$$

The nozzle boundary is defined by $v_0 = \text{constant}$, which along with the equation for the nozzle wall in cartesian

coordinates, $y' = f(x)$, implicitly defines the constant a . Substituting for μ in the velocity potential produces,

$$= A \cosh^{-1} [(r_1 + r_2)/a]$$

from which the velocities may be determined.

$$u = \phi_x = \frac{A}{\sqrt{[(r_1 + r_2)/a]^2 - 1}} \left\{ a \frac{\partial r_1}{\partial x} + a \frac{\partial r_2}{\partial y} \right\}$$

$$v = \phi_y = \frac{A}{\sqrt{[(r_1 + r_2)/a]^2 - 1}} \left\{ a \frac{\partial r_1}{\partial y} + a \frac{\partial r_2}{\partial x} \right\}$$

$$\frac{\partial r_1}{\partial x} = \frac{x}{\sqrt{(y + a/2)^2 + x^2}}$$

$$\frac{\partial r_2}{\partial x} = \frac{x}{\sqrt{(y - a/2)^2 + x^2}}$$

$$\frac{\partial r_1}{\partial y} = \frac{y + a/2}{\sqrt{(y + a/2)^2 + x^2}}$$

$$\frac{\partial r_2}{\partial y} = \frac{y - a/2}{\sqrt{(y - a/2)^2 + x^2}}$$

The constant of integration A may be determined by specifying the flow rate through the nozzle, but when the inlet velocities are normalized with respect to the freestream velocity (U_∞) A is factored out of the problem.

Velocities for compressible flow can be solved by mapping back to the physical coordinate system (x, y plane).

Other boundary conditions are universal to both problems. These are:

$$u = 0 \quad \text{on the line of symmetry}$$

$$v = (1 + u)df/dx \quad \text{at the nozzle wall}$$

$F(x)$ defines the nozzle boundary in terms of a ratio of the throat semi-height as a function of x . The throat semi-height is taken to be 1 for convenience.

Pressure ratio, sound speed, and Mach number are calculated as before by equations 5 through 8.

VIII. DISCUSSION OF RESULTS

A. TIME INTEGRATION TO A STEADY STATE SOLUTION

As stated before, the Houbolt method of integration is stable for all time steps. Results of the test cases bear this out with the larger time steps providing the most rapid convergence to a steady state solution. Time steps were tried from $t = .1$ to $t = 100$. Time steps larger than $t = 100$ were not attempted because as t becomes too large the influence that the damping and mass matrices have on the effective stiffness matrix becomes negligible compared to the stiffness matrix. That is:

$$SE_{ij} = S_{ij} + 2SC_{ij}/\Delta t^2 + 6SM_{ij}/\Delta t$$

as $\Delta t \rightarrow \infty$

$$SE_{ij} \approx S_{ij}$$

The starting solutions were chosen somewhat arbitrarily. $\phi_{j,2\Delta t}$ was chosen to satisfy the first iteration of the steady solution

$$S_{ij}\phi_{j,2\Delta t} = 0$$

when the non-linear term (u) in the coefficient

$$1 - M_{\infty}^2 - M_{\infty}^2(\gamma + 1) u$$

was set to zero. $\phi_{j,\Delta t}$ and $\phi_{j,0}$ were chosen as multiples of $\phi_{j,2\Delta t}$ and respectively they were

$$\phi_{j,\Delta t} = .5\phi_{j,2\Delta t}$$

$$\phi_{j,0} = 0$$

This starting procedure proved to be superior to choosing the first three vectors closer to the converged solution. If $\phi_{j,2\Delta t}$, $\phi_{j,\Delta t}$, and $\phi_{j,0}$ were chosen to be the last three time steps of the previous case, the solution oscillated and converged much slower than with the starting solutions chosen as above.

The stiffness, mass, and damping matrices were recalculated after each time step, using the under-relaxation technique described above. This was necessary to utilize the special assembly procedures invoked by STRANL-II to prevent the inadmissible influence of downwind nodes from propagating upstream in the supersonic region.

For barely critical flow ($M_\infty = .861$) and subsonic flow, an under-relaxation factor $\theta = 1$ (successive approximation) resulted in convergence to a steady state solution after only three time steps. Eleven time steps were required for the supercritical solution to converge using the same relaxation factor. Reducing θ to .5 increased the rate of convergence and the solution achieved steady state after six time steps. Figure 4 compares the steady state solution for a 6% thick circular arc airfoil at $M_\infty = .909$, using the same integration method, with the results obtained in Ref. [5]. Chan's results converged in 10 iterations after using the results from the barely critical flow as an initial guess to the

supercritical solution. Figure 4 is a plot of local Mach numbers at boundary nodes on the airfoil.

B. CONVERGING-DIVERGING NOZZLE

The nozzle chosen for the test cases was the two-dimensional Oswatitsch nozzle with the boundary defined as

$$y = 1 + \sqrt{.2(x - 2.5)^2}$$

where the throat at $x = 2.5$ has a semi-height of 1. The inlet was taken to be $x = 0$ and the exit was at $x = 5$. $M_\infty = .44$, the inlet Mach number on the nozzle center-line was chosen to yield sonic conditions in the throat.

Two solutions were possible for this inlet condition,- the symmetric solution and the asymmetric solution; but neither solution was achieved by the finite element method. Although the solution converged for the subsonic case in three iterations, center-line Mach numbers deviated significantly from both one-dimensional theory and from Oswatitsch's approximation [Fig. 5]. When the local Mach number M exceeded the inlet Mach number by approximately .2 ($M \geq .64$) the solution was invalid. Differences at the center part of the nozzle are due to an essentially incorrect free stream Mach number. Patching the solution at $x \approx 1.5$ would improve the solution.

A second test case was run for the supersonic section of the nozzle with the inlet boundary on the sonic line. The exit boundary was left free to float. Here the solution was unstable and no meaningful results were obtained.

— Steady solution
+ Time integration

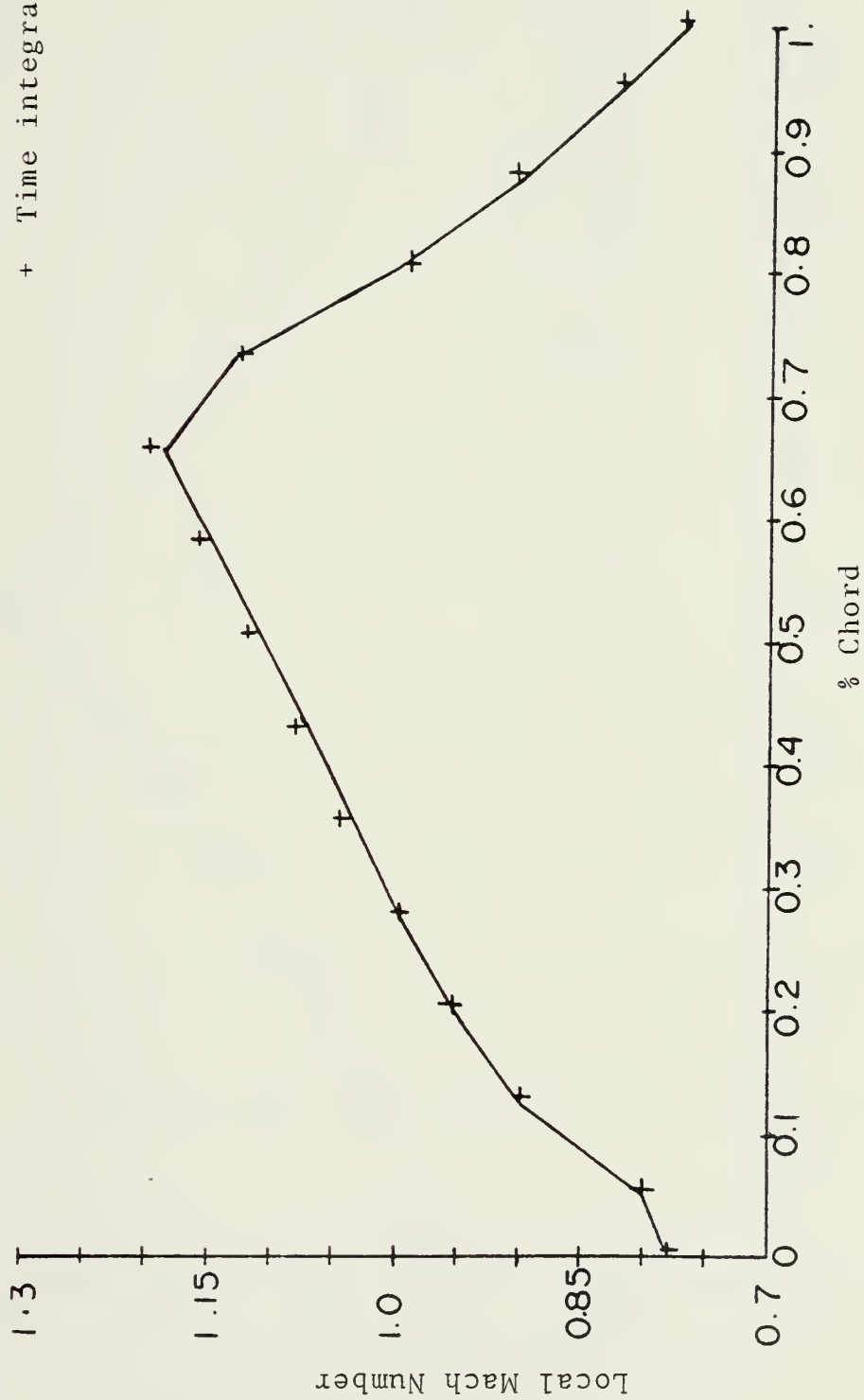


Figure 4 - Comparison of Time Integration Results with Steady State Results.

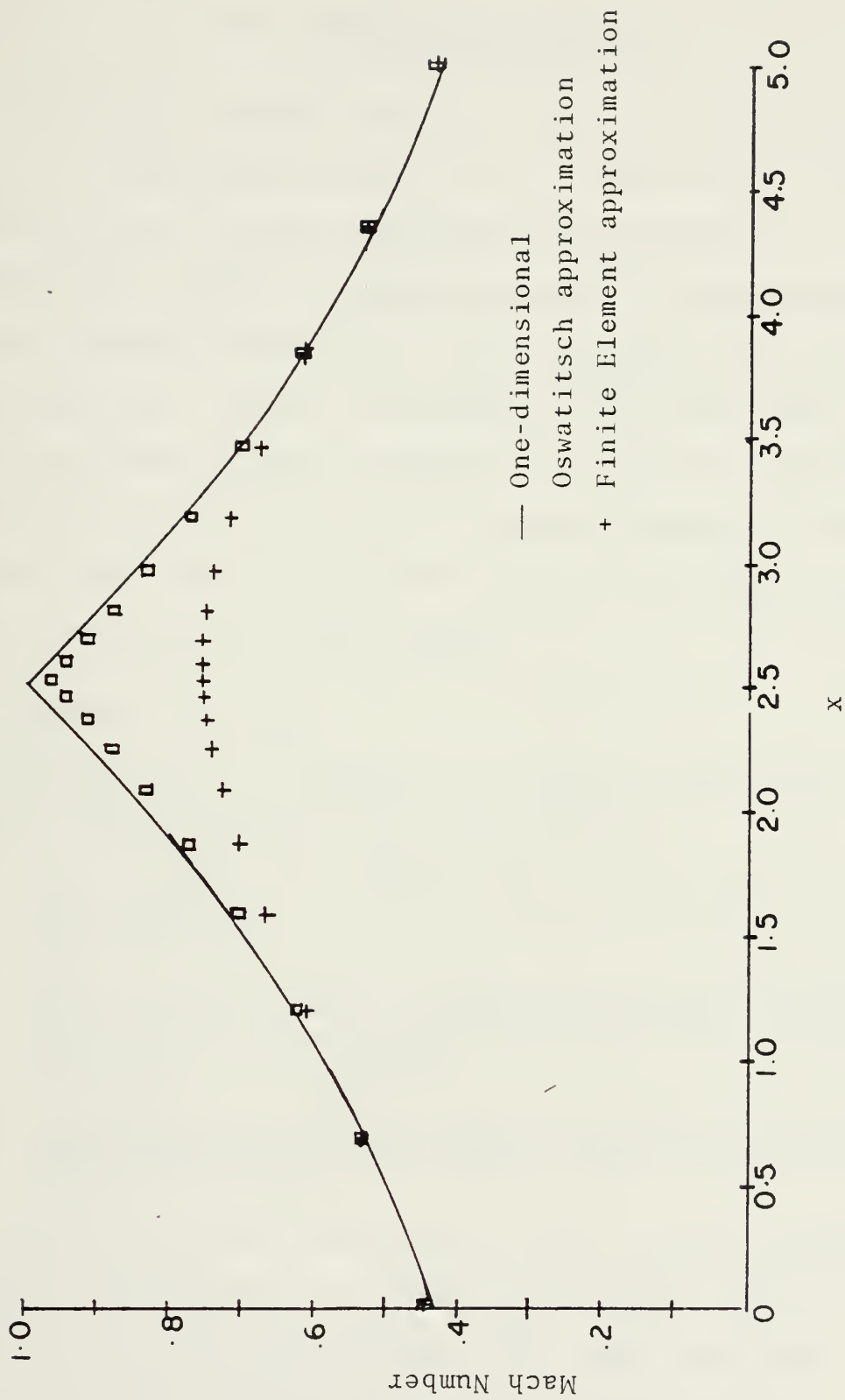


Figure 5 - Nozzle Center-Line, Inlet Mach Number 0.4300.

IX. PROGRAM MODIFICATION

The finite element computer program for non-lifting air-voils, as developed in Ref. [5], is separated into two parts. These have been designated STRANL-I and STRANL-II by Lockheed Corporation. STRANL-I generates a finite element mesh to be used as inputs to STRANL-II, which assembles the finite element equations, applies the boundary conditions, and solves the non-linear system of equations. Detailed descriptions and instructions for the use of the two programs can be found in Ref. [5]. Only the modifications to the above programs will be discussed in this section.

A. UNSTEADY EQUATIONS

Modifications to STRANL-II to form and solve the unsteady finite element equations were three-fold:

- 1) The new elemental matrices SC_{ij} and SM_{ij} were calculated and assembled.
- 2) All the matrices were stored on an external magnetic disk to be accessed and reassembled later because of the amount of space required to store three large matrices, in core memory.
- 3) The effective stiffness matrix SE_{ij} and the effective load vector RE_{ij} were assembled, and the system of equations solved.

Several existing subroutines in the original STRANL-II program were modified to assemble the damping and mass matrices. These include subroutines NEWK, EMTC, DERV, and EMQT. Two new subroutines were added to perform the other tasks.

B. MODIFICATIONS TO CALCULATE THE MASS AND DAMPING MATRICES

EMTC in the STRANL-II program calculated the elemental stiffness matrix by numerically integrating the equation,

$$S_{ij} = Q_i P_j dA$$

Equations were added to EMTC to perform the additional numerical integrations for the mass and damping matrices. All three matrices were calculated at the same time. EMQT assembled the elemental stiffness matrices for a quadrilateral and trapezoidal element from the contributions of the triangular elements. Mass and damping matrices were calculated in the same fashion.

Subroutine NEWK, an original subroutine in STRANL-II, which assembled the finite element equations for steady flow, was modified to assemble SC_{ij} and SM_{ij} . A new calling argument, NMAT was passed to NEWK, which assembled contributions from the triangular, quadrilateral and trapezoidal element matrices into the global matrices S_{ij} , SC_{ij} , and SM_{ij} , depending on NMAT being 1, 2 or 3.

1. Subroutine STORE

Given a non-symmetric matrix stored in a banded node, plus the right hand side vector, subroutine STORE separates this system into two matrices and stores them on a magnetic disk. Figures 6 and 7 show the decomposition of a banded matrix into banded storage, and the further decomposition of this banded stored matrix to two smaller matrices by subroutine STORE. In these figures, D, L, and U represent the

$N \times N$

$N \times 2HBW$

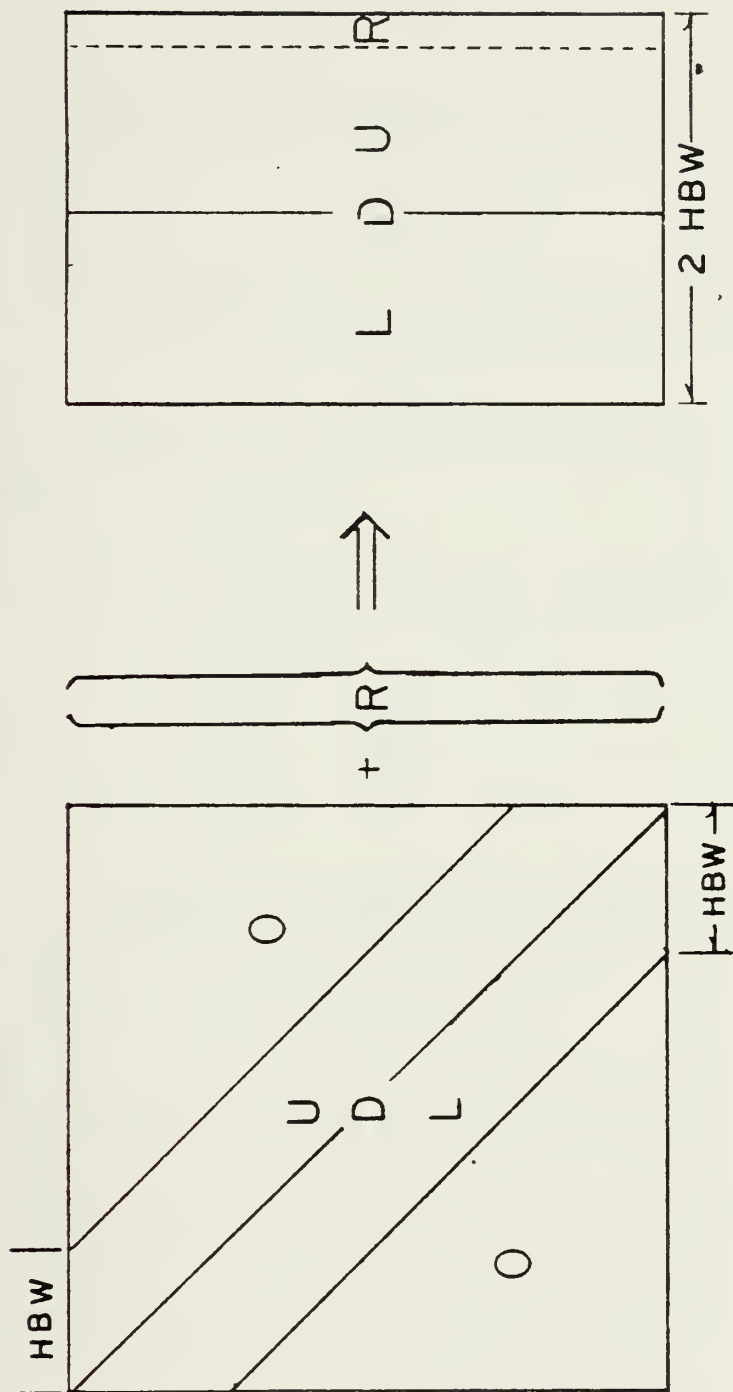


Figure 6 - Decomposition of a Banded Matrix Plus Right Hand Side

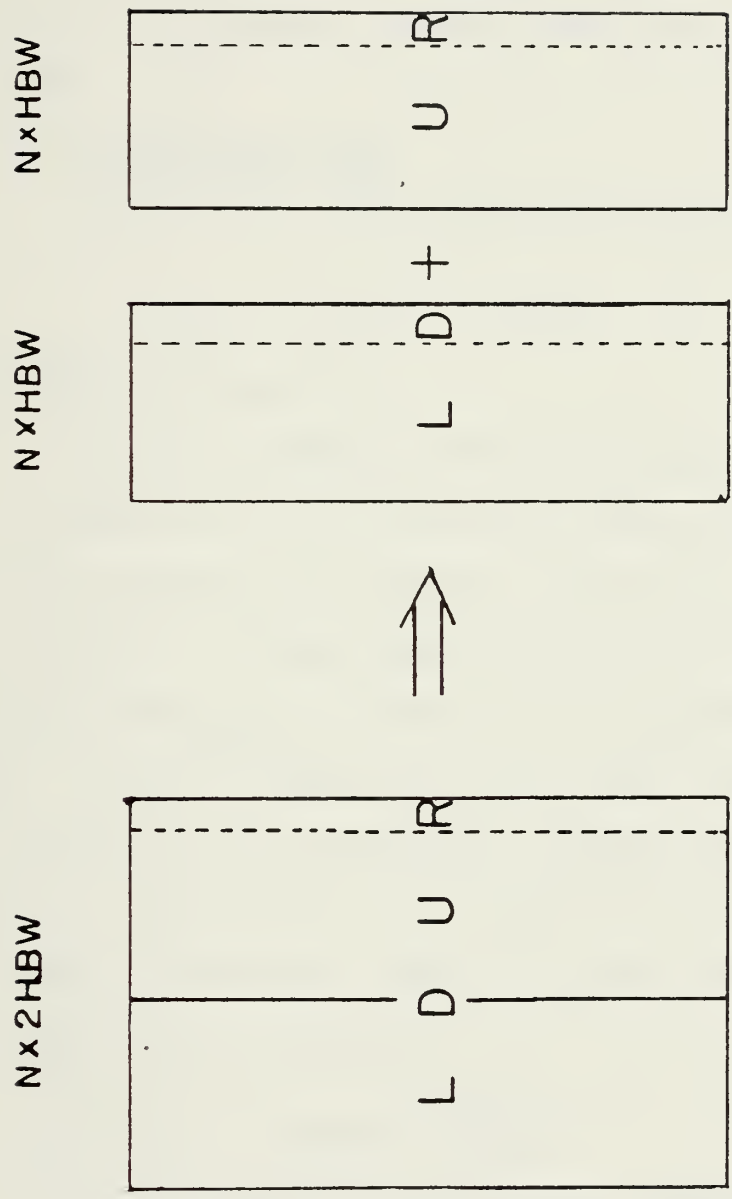


Figure 7 - Separation of a Banded Stored Matrix by Store

diagonal matrix, the lower triangular matrix, and the upper triangular matrix respectively. HBW is the half bandwidth and R is the right hand side vector.

STORE requires an additional work area one-half the size of the originally dimensioned matrix which is to be stored.

2. Subroutine TIME

Subroutine TIME integrates the system

$$S_{ij}\phi_j + SC_{ij}\dot{\phi}_j + SM_{ij}\ddot{\phi}_j = R_j$$

by Houbolt integration.

TIME reassembles the three matrices which were stored on the magnetic disk to form the effective stiffness matrix and the effective load vector. Once this system of equations is assembled, a banded equation solver is called to yield the solution for $\phi_{j,t+\Delta t}$. Figure 8 is a schematic flow chart of TIME. In Fig. 8 when $L = 1$, the lower triangular matrix and the diagonal of the effective stiffness matrix are formed by adding the appropriate contributions from the stiffness, mass and damping matrices. When $L = 2$, the upper triangular matrix is formed in like fashion.

C. CONVERGING-DIVERGING NOZZLE

1. Application of the Boundary Conditions

Regardless of the type of problem for which a set of system equations have been assembled, the equations will have the form

$$K_{ij} x_i = R_i$$

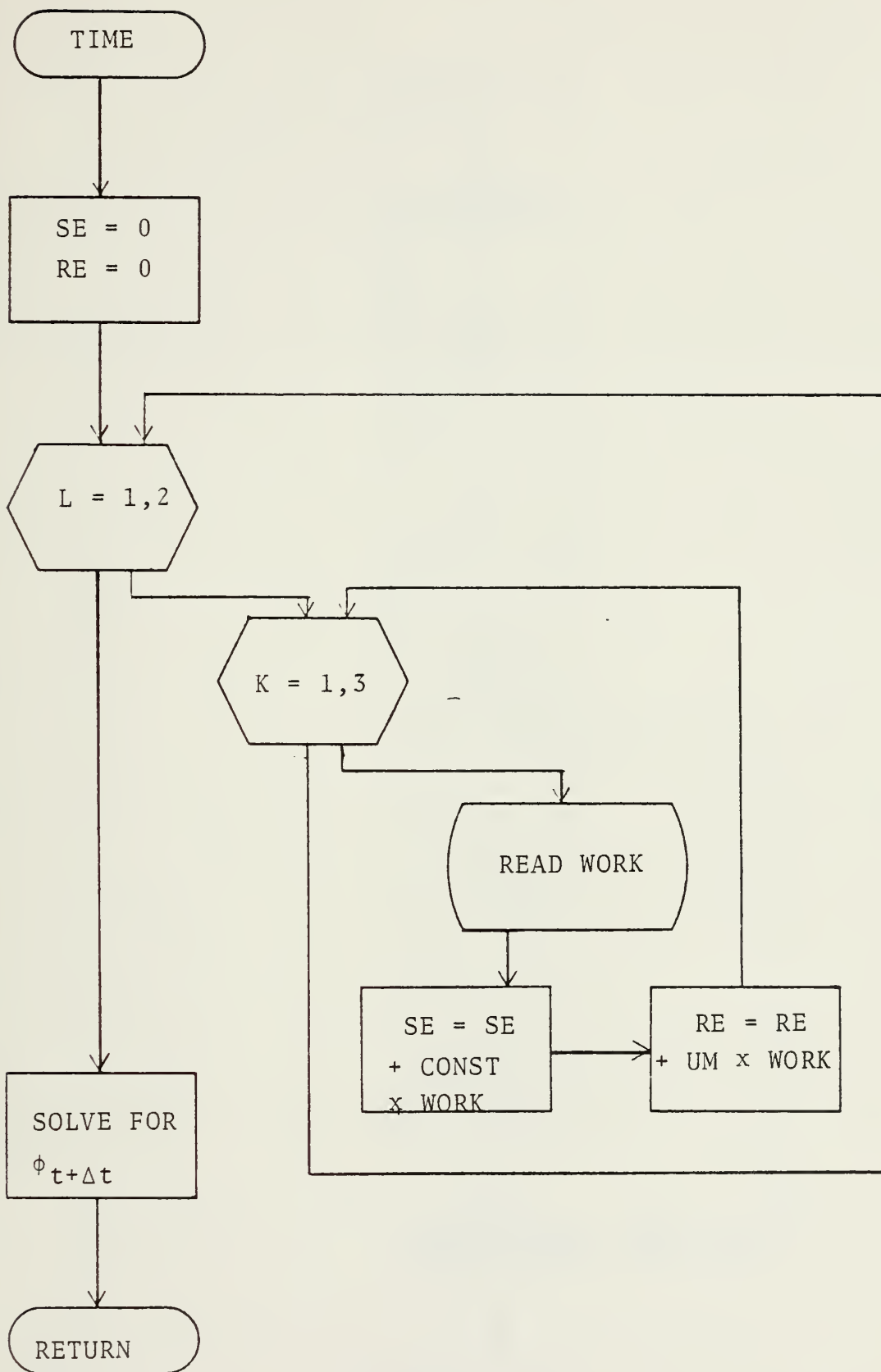


Figure 8 - Flow Chart of Subroutine TIME

STRANL-II

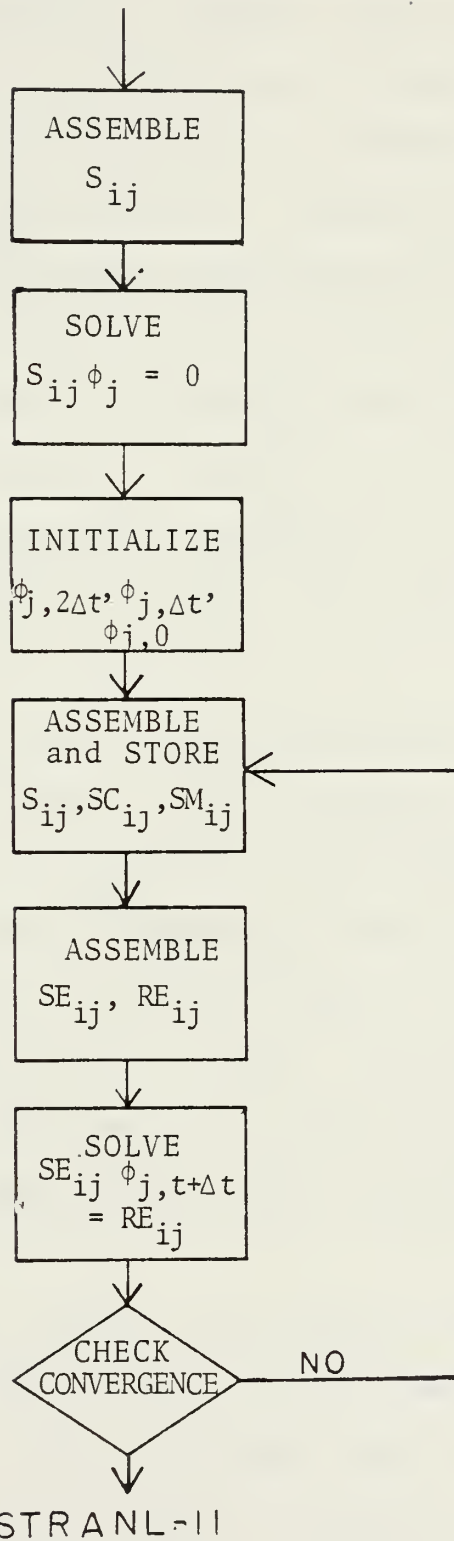


Figure 9 - Flow Chart of STANL-II Modification to Integrate Unsteady Equations

in which K_{ij} is an $n \times n$ matrix and x_i and R_i are vectors of length n . These equations do not take into account the known values of x_i on the boundaries. However, for a unique solution of the above equation, at least one or more nodal variables must be specified and K_{ij} must be modified to render it non-singular. For each equation i , either x_i or R_i must be specified but it is physically impossible to specify both x_i and R_i . There are a number of ways to apply the boundary conditions to the equations and when they are applied the number of equations is reduced. However, it is convenient to leave the number of equations unchanged to avoid major restructuring of the computer storage. One such method is described below.

If k is the subscript of the prescribed nodal variable, the k^{th} row and the k^{th} column of the original K_{ij} matrix are set to zero, K_{kk} is set to 1 and R_k is replaced by the known value of x_k . Each of the $n-1$ remaining terms of R_i is modified by subtracting from it the value $K_{ik}x_k$. This procedure is repeated for all the boundary values. Of course, when the matrix is stored in a banded mode, the algorithm will differ from that for the $n \times n$ square matrix, but the procedure is similar.

Subroutine BNDRY applies the boundary conditions for the modified program. Setting the option parameter IOPT(4), in STRANL-II, equal to 1 will call BNDRY which will read the boundary velocities and apply them to a banded stored matrix by the method described above.

When the value of the x_k is zero, as in the non-lifting airfoil problem, the algorithm becomes simpler than the above method because there is no need to either set the k^{th} column to zero or subtract the value $K_{ik}x_k$ from the right hand vector.

X. TEST CASES

In computing the flow field for either the non-lifting airfoil or for the Laval nozzle, the following procedures were followed:

- 1) The desired mesh was sketched with each node assigned a number.
- 2) Appropriate input cards based on the sketch were prepared and supplied to STRANL-I to generate the data on punched cards for input to STRANL-II.
- 3) The above punched cards were supplied to STRANL-II with three additional cards as input parameters for each case, plus additional cards for the boundary velocities, if the nozzle solution is desired.
- 4) Results of the finite element calculations are printed after each iteration, and the converged solution is punched on cards for possible later use.

A. TIME INTEGRATION TO A STEADY STATE SOLUTION

Test cases for the integration of the unsteady transonic finite element equations were conducted to calculate the steady transonic flow over a 6% thick circular arc airfoil. These tests were made using the same airfoil, mesh, and free-stream Mach numbers as Chan et al. [Ref. 5] published. These conditions were chosen to provide a source for comparison of the results.

Freestream Mach numbers used in these calculations were:

$$M_{\infty} = .806 \quad (\text{subcritical})$$

$$M_{\infty} = .861 \quad (\text{barely critical})$$

$$M_{\infty} = .909 \quad (\text{supercritical}).$$

Each case was treated individually with $\phi_j = 0$ used as the initial guess for each case, whereas Chan et al. [Ref. 5] used zero for the initial guess for $M_\infty = .806$ and then used the computed results as the initial guess for each subsequent case.

DELT, the value for the time step, is input by a parameter specified in columns 41-45 of the second card following the title card for each case when the unsteady option (IOPT(6) = 1) is selected.

1. STRANL-I Program

a. Input

Input cards used to generate the finite element mesh are listed on the next page. Cards were arranged in accordance with Ref. [5], in the following order:

- Title card
- Option card
- Element cards
- Card for the total number of nodes
- Node coordinate cards
- Card for the number of boundary nodes
- Card for the boundary nodes at infinity
- Card for the nodes on the line of symmetry
- Card for the nodes on the airfoil
- Cards for the slope of the airfoil.

Input cards to STRANL-I for these tests are listed on the next three pages.

b. Output

Output from STRANL-I is in the form of printouts and punched cards. Printouts from STRANL-I are listed on the following eight pages.

INPUT TO STRANL-I

STEADY	TRANSCNIC	FLOW	MESH	6--154	ELEMENTS,	170	NODES	
3	4	1	0	31	21	26	TRA000010	
4	1	0	1	6	1	22	TRA000020	
3	1	0	1	27	28	33	TRA000030	
4	1	0	0	33	34	36	TRA000040	
3	1	0	0	33	34	37	TRA000050	
4	1	0	3	48	49	53	TRA000060	
4	1	0	0	51	52	55	TRA000070	
3	1	0	0	55	56	57	TRA000080	
4	1	0	0	55	56	60	TRA000090	
3	1	0	0	55	56	62	TRA000100	
4	1	0	0	55	56	67	TRA000110	
3	1	0	0	55	56	67	TRA000120	
4	1	0	5	61	62	67	TRA000130	
4	1	0	0	65	66	67	TRA000140	
4	1	0	0	10	11	67	TRA000150	
4	1	0	0	15	16	67	TRA000160	
4	1	0	0	20	21	67	TRA000170	
4	1	0	0	25	26	67	TRA000180	
4	1	0	0	30	31	67	TRA000190	
4	1	0	0	35	36	67	TRA000200	
4	1	0	0	38	39	67	TRA000210	
4	1	0	0	41	42	67	TRA000220	
4	1	0	0	44	45	67	TRA000230	
4	1	0	0	47	48	67	TRA000240	
4	1	0	0	47	48	67	TRA000250	
4	1	0	0	50	51	67	TRA000260	
4	1	0	0	52	53	67	TRA000270	
4	1	0	0	58	59	67	TRA000280	
4	1	0	0	63	64	67	TRA000290	
4	1	0	0	67	68	67	TRA000300	
4	1	0	0	70	71	67	TRA000310	
4	1	0	0	75	76	67	TRA000320	
4	1	0	0	80	81	67	TRA000330	
4	1	0	0	85	86	67	TRA000340	
3	1	0	0	85	86	67	TRA000350	
4	1	0	0	92	93	67	TRA000360	
4	1	0	0	94	95	67	TRA000370	
4	1	0	0	96	97	67	TRA000380	
3	1	0	0	99	100	67	TRA000390	
4	1	0	0	102	103	67	TRA000400	
4	1	0	0	104	105	67	TRA000410	
4	1	0	0	109	110	67	TRA000420	
4	1	0	0	109	110	67	TRA000430	
4	1	0	0	146	147	67	TRA000440	
4	1	0	0	151	152	67	TRA000450	
4	1	0	0	151	152	67	TRA000460	
4	1	0	0	158	159	67	TRA000470	
4	1	0	0	158	159	67	TRA000480	
4	1	0	0	158	159	67	TRA000490	
4	1	0	0	158	159	67	TRA000500	

86	89	93	98	104	111	118	125	132	135	146	153	159	164	168
.06022	+0.10820	+0.09004	+0.09004	+0.09004	+0.07193	+0.07193	+0.05385	+0.05385	+0.03589	+0.03589	+0.01794	+0.01794	+0.00000	+0.00000
-.01794	-0.03589	-0.05388	-0.05388	-0.05388	-0.07193	-0.07193	-0.09004	-0.09004	-0.10820	-0.10820	-0.01794	-0.01794	-0.00000	-0.00000
1.0	8.0	0.1	0.1	0.1	3.5	3.5	0.1	0.1	0.1	0.1	0.1	0.1	0.1	0.1

CIRCULAR ARC

[illegible]

98	0.24250E	01
97	0.24250E	01
96	0.24250E	01
95	0.24250E	01
94	0.24250E	01
90	0.25000E	01
89	0.25000E	01
88	0.25000E	01
87	0.25000E	01
86	0.25000E	01
85	0.25000E	01
84	0.25000E	01
80	0.25750E	01
79	0.25750E	01
78	0.25750E	01
77	0.25750E	01
76	0.25750E	01
75	0.25750E	01
74	0.25750E	01
70	0.26500E	01
69	0.26500E	01
68	0.26500E	01
67	0.26500E	01
66	0.26500E	01
65	0.26500E	01
64	0.26500E	01
60	0.27250E	01
59	0.27250E	01
58	0.27250E	01
57	0.27250E	01
56	0.27250E	01
55	0.27250E	01
54	0.27250E	01
49	0.28000E	01
48	0.28000E	01
47	0.28000E	01
46	0.28000E	01
45	0.28000E	01
44	0.28000E	01
37	0.28750E	01
36	0.28750E	01
35	0.28750E	01
34	0.28750E	01
33	0.28750E	01
27	0.29500E	01
26	0.29500E	01
25	0.29500E	01
24	0.29500E	01

2. STRANL-II Program

a. Input

Listed on the next three pages are the input cards to the STANL-II program.

b. Output

The output from this program is in the form of printouts for each iteration and punched cards for the converged solution. Output from STRANL-II for $M_\infty = .909$ is listed on the following eight pages.

[illegible]

111
161 3150
14097
187 6536
133
173 303
14108
1186637
126
172 292
14259
1187744
125
171 289
14350
116845
124
183 300
14974
1159946
123
182 320
15085
118707
122
181 420
15156
1197483
138
193 418
1520
119754914
137
192 407
1581
119876415
136
191 388
15979
112775524
135
103 330
16028
11005625
134
102 2216
1294
11045726
148
101 191
1305
1105805827
147
114 174
1316
110645933
146
113 36
1112
110756034
145
112 2225
13586
110864535

TIME INTEGRATION TO STEADY SOLUTION -- CIRCULAR ARC --170 NODES--M=0.909
CONVERGENCE LIMIT =0.0050

C C C C C 1 C C C C C C C C C C C C C C C C

NO. OF ELEMENTS= 154 NO. OF NODES= 170 FULL BANDWIDTH = 84

ELE. NO. ARC ELEMENT NCCES

1	31	21	26	C
2	11	11	11	11
3	11	11	11	11
4	11	11	11	11
5	11	11	11	11
6	11	11	11	11
7	11	11	11	11
8	11	11	11	11
9	11	11	11	11
10	11	11	11	11
11	11	11	11	11
12	11	11	11	11
13	11	11	11	11
14	11	11	11	11
15	11	11	11	11
16	11	11	11	11
17	11	11	11	11
18	11	11	11	11
19	11	11	11	11
20	11	11	11	11
21	11	11	11	11
22	11	11	11	11
23	11	11	11	11
24	11	11	11	11
25	11	11	11	11
26	11	11	11	11
27	11	11	11	11
28	11	11	11	11
29	11	11	11	11
30	11	11	11	11
31	11	11	11	11
32	11	11	11	11
33	11	11	11	11
34	11	11	11	11
35	11	11	11	11
36	11	11	11	11
37	11	11	11	11
38	11	11	11	11
39	11	11	11	11
40	11	11	11	11
41	11	11	11	11
42	11	11	11	11
43	11	11	11	11
44	11	11	11	11
45	11	11	11	11
46	11	11	11	11
47	11	11	11	11
48	11	11	11	11
49	11	11	11	11
50	11	11	11	11
51	11	11	11	11
52	11	11	11	11
53	11	11	11	11
54	11	11	11	11
55	11	11	11	11
56	11	11	11	11
57	11	11	11	11
58	11	11	11	11
59	11	11	11	11
60	11	11	11	11
61	11	11	11	11
62	11	11	11	11
63	11	11	11	11
64	11	11	11	11
65	11	11	11	11
66	11	11	11	11
67	11	11	11	11
68	11	11	11	11
69	11	11	11	11
70	11	11	11	11
71	11	11	11	11
72	11	11	11	11
73	11	11	11	11
74	11	11	11	11
75	11	11	11	11
76	11	11	11	11
77	11	11	11	11
78	11	11	11	11
79	11	11	11	11
80	11	11	11	11
81	11	11	11	11
82	11	11	11	11
83	11	11	11	11
84	11	11	11	11

139	70	0.26500E	01	0.27310E	01
140	69	0.26500E	01	0.12120E	00
141	68	0.26500E	01	0.23400E	00
142	67	0.26500E	01	0.36920E	00
143	66	0.26500E	C1	0.33150E	00
144	65	0.26500E	01	0.72630E	00
145	64	0.26500E	01	0.96000E	00
146	60	0.27250E	C1	0.23940E	01
147	59	0.27250E	01	0.11420E	00
148	58	0.27250E	C1	0.22250E	00
149	57	0.27250E	C1	0.35240E	00
150	56	0.27250E	01	0.50830E	00
151	55	0.27250E	C1	0.69550E	00
152	54	0.27250E	C1	0.92000E	00
153	49	0.28000E	C1	0.19220E	01
154	48	0.28000E	C1	0.11740E	00
155	47	0.28000E	01	0.23530E	00
156	46	0.28000E	01	0.37670E	00
157	45	0.28000E	01	0.54640E	00
158	44	0.28000E	01	0.75000E	00
159	37	0.28750E	C1	0.13150E	01
160	36	0.28750E	C1	0.11320E	00
161	35	0.28750E	C1	0.23320E	00
162	34	0.28750E	C1	0.37720E	00
163	33	0.28750E	01	0.55000E	00
164	27	0.29500E	01	0.57230E	02
165	26	0.29500E	C1	0.10030E	00
166	25	0.29500E	01	0.21380E	CG
167	24	0.29500E	01	0.35000E	00
168	15	0.30000E	01	0.0	00
169	14	0.30000E	C1	0.10000E	00
170	13	0.30000E	C1	0.22000E	00

$\text{NACH NUMBER} = 0.569 \quad \text{RELAX. FACTOR} = 0.5000$

AC. ITERATIONS = 7

NCLE	PHI I	UCUM	UCUM	CLT	LMAC	P/P/C	CP	DELM
1	1.704E-03	1.000E-00	1.905E-06	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
2	1.414E-04	1.001E-00	1.469E-06	8.796E-01	5.024E-01	5.842E-01	1.192E-05	0.181E-03
3	1.313E-04	1.540E-01	1.901E-05	7.147E-01	5.024E-01	5.842E-01	1.192E-05	0.332E-03
4	1.359E-05	9.243E-01	5.557E-04	7.761E-01	5.326E-01	5.383E-01	1.117E-02	1.543E-03
5	1.309E-05	1.000E-00	5.721E-04	8.284E-01	5.101E-01	5.850E-01	1.211E-03	0.0
6	1.434E-05	1.001E-00	5.571E-04	8.284E-01	5.101E-01	5.850E-01	1.211E-03	0.303E-03
7	1.402E-05	9.815E-01	5.547E-04	7.857E-01	5.033E-01	5.850E-01	1.063E-02	1.206E-04
8	1.770E-05	9.287E-01	1.000E-00	7.857E-01	5.033E-01	5.850E-01	1.063E-02	1.620E-04
9	1.204E-05	1.000E-00	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
10	1.784E-05	1.000E-00	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
11	1.384E-05	9.963E-01	8.874E-04	8.284E-01	5.033E-01	5.850E-01	1.063E-02	0.0
12	1.304E-05	9.963E-01	8.874E-04	8.284E-01	5.033E-01	5.850E-01	1.063E-02	0.0
13	1.304E-05	9.963E-01	8.874E-04	8.284E-01	5.033E-01	5.850E-01	1.063E-02	0.0
14	1.304E-05	9.963E-01	8.874E-04	8.284E-01	5.033E-01	5.850E-01	1.063E-02	0.0
15	1.304E-05	9.963E-01	8.874E-04	8.284E-01	5.033E-01	5.850E-01	1.063E-02	0.0
16	1.304E-05	9.963E-01	8.874E-04	8.284E-01	5.033E-01	5.850E-01	1.063E-02	0.0
17	1.304E-05	9.963E-01	8.874E-04	8.284E-01	5.033E-01	5.850E-01	1.063E-02	0.0
18	1.304E-05	9.963E-01	8.874E-04	8.284E-01	5.033E-01	5.850E-01	1.063E-02	0.0
19	1.304E-05	9.963E-01	8.874E-04	8.284E-01	5.033E-01	5.850E-01	1.063E-02	0.0
20	1.304E-05	9.963E-01	8.874E-04	8.284E-01	5.033E-01	5.850E-01	1.063E-02	0.0
21	1.304E-05	9.963E-01	8.874E-04	8.284E-01	5.033E-01	5.850E-01	1.063E-02	0.0
22	1.304E-05	9.963E-01	8.874E-04	8.284E-01	5.033E-01	5.850E-01	1.063E-02	0.0
23	1.304E-05	9.963E-01	8.874E-04	8.284E-01	5.033E-01	5.850E-01	1.063E-02	0.0
24	1.304E-05	9.963E-01	8.874E-04	8.284E-01	5.033E-01	5.850E-01	1.063E-02	0.0
25	1.304E-05	9.963E-01	8.874E-04	8.284E-01	5.033E-01	5.850E-01	1.063E-02	0.0
26	1.304E-05	9.963E-01	8.874E-04	8.284E-01	5.033E-01	5.850E-01	1.063E-02	0.0
27	1.304E-05	9.963E-01	8.874E-04	8.284E-01	5.033E-01	5.850E-01	1.063E-02	0.0
28	1.304E-05	9.963E-01	8.874E-04	8.284E-01	5.033E-01	5.850E-01	1.063E-02	0.0
29	1.304E-05	9.963E-01	8.874E-04	8.284E-01	5.033E-01	5.850E-01	1.063E-02	0.0
30	1.304E-05	9.963E-01	8.874E-04	8.284E-01	5.033E-01	5.850E-01	1.063E-02	0.0
31	1.304E-05	9.963E-01	8.874E-04	8.284E-01	5.033E-01	5.850E-01	1.063E-02	0.0
32	1.304E-05	9.963E-01	8.874E-04	8.284E-01	5.033E-01	5.850E-01	1.063E-02	0.0
33	1.304E-05	9.963E-01	8.874E-04	8.284E-01	5.033E-01	5.850E-01	1.063E-02	0.0
34	1.304E-05	9.963E-01	8.874E-04	8.284E-01	5.033E-01	5.850E-01		

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B. CONVERGING-DIVERGING NOZZLE

Two test cases were run for the converging-diverging nozzle. The first case was for symmetric flow designed to yield sonic conditions in the throat of the nozzle. The second case dealt with supersonic flow in the diverging section by starting with the sonic line as the boundary of the nozzle mesh. Oswatitsch's two-dimensional nozzle [Ref. 9], $y = 1 + \sqrt{2(x - 2.5)^2}$ with a semi-throat height of 1 at $x = 2.5$ was used in both cases. Boundaries for the subsonic nozzle were at $x = 0$ and $x = 5$.

1. Symmetric Solution

This problem was analyzed using the mesh shown in Fig. 10, which consists of 126 elements and 152 nodes. In the second card (the option card) IOPT(4) = 1 indicates that non-zero boundary velocities at the inlet and the exit will be read and applied by subroutine BNDRY. This option requires that the number of inlet and exit boundary nodes be specified in columns 36-40 of the next card. Perturbation velocities at the boundary nodes follow on the subsequent four cards. Subroutine BNDRY reads u and v respectively for the first boundary node and then continues reading u and v for each inlet and then each exit node in the order specified on the appropriate card.

2. Supersonic Case--Diverging Section

The mesh used for this case is sketched in Fig. 11 and input cards to STRANL-II follow on the next page. For this

case the options in effect are IOPT(4) = 1 and IOPT(5) = 1 which cause non-zero boundary velocities to be applied to the sonic line only.

STEADY TRANSONIC FLOW CONVERGING DIVERGING NOZZLE-3 M=.375

[illegible]

113	114	115	116	117	118	119	120	121	122	123	124	125	126	127	128
129	130	131	132	133	134	135	136	137	138	139	140	141	142	143	144
145	146	147	148	149	150	151	152								
STEADY FLOW-CONVERGING DIVERGING NOZZLE M=.4388															
1	1														
1C	.01		.4388			1.0	16								
0.0	0.0		-8.571E-04	-3.995E-02	-3.473E-03	-8.004E-02	-7.978E-03	-1.204E-01							
-1.458E-C2	-1.611E-01	-2.356E-02	-2.021E-01	-3.525E-01	-4.997E-02	-2.846E-C1									
0.0	0.0		-8.571E-04	3.995E-02	-3.473E-03	8.004E-02	-7.978E-03	1.204E-01							
-1.458E-C2	1.611E-01	-2.356E-02	2.021E-01	-3.525E-02	2.433E-01	-4.997E-02	2.846E-01								

STEADY SUPersonic FLOW, SONIC LINE INPUT, UNIFORM MESH

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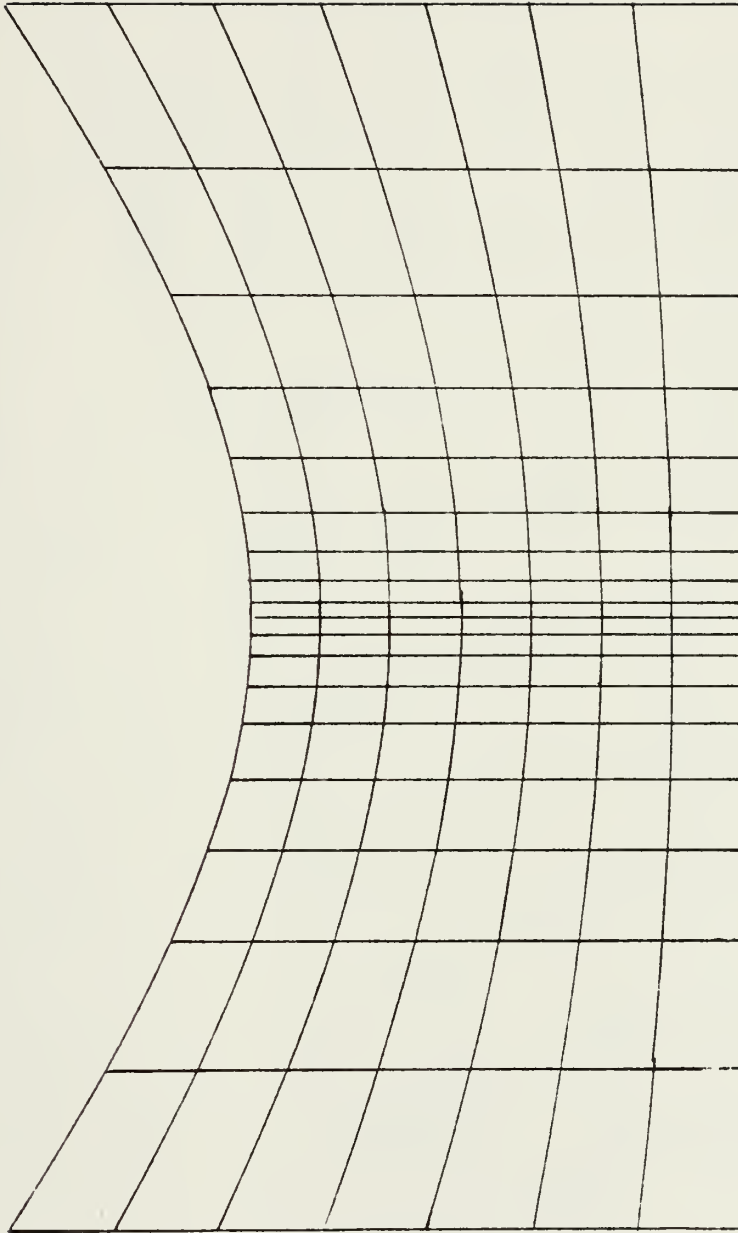


Figure 10 - Finite Element Mesh for the Converging-Diverging Nozzle.

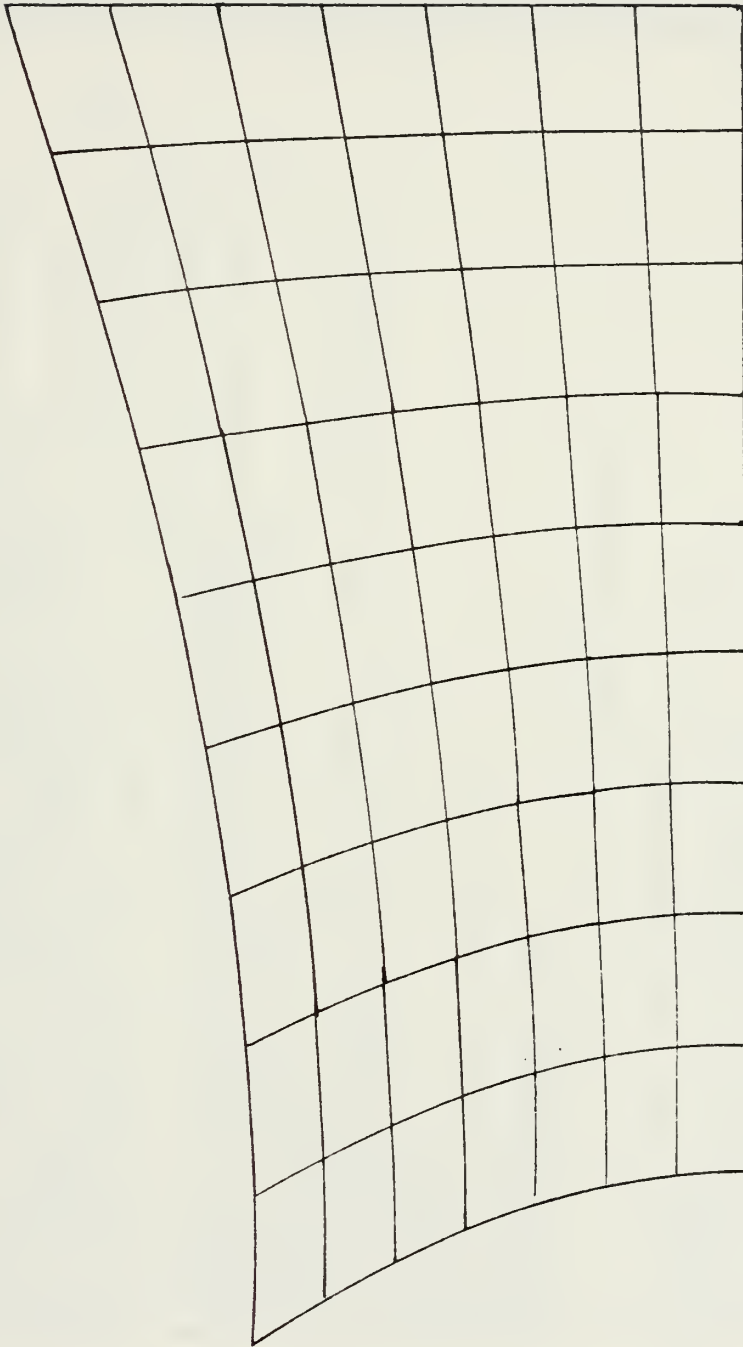


Figure 11 - Finite Element Mesh for the Supersonic Case


```

C
C
C
      IF (IOPT(4) .EQ. 1) READ(NREAD,840)(UB(I),VB(I),I=1,NFARF$)
      IF (IOPT(4) .EQ. 1) CALL BCCND(UB,VB,IOPT(5),NFARF$)
      IF (IOPT(1) .EQ. 1) GO TO 382

      READ AND PRINT MESH DATA, BOUNDARY NODES, AND AIRFCIL SLOPE

      READ (NREAD,825) NELS,NPS,NBW,(NICS(I),I=1,3)
      READ (NREAD,825) ((NOD(I,J),J=1,4),I=1,NELS)
      READ (NREAD,840) (X(I),Y(I),I=1,NPS)
      CC 110 I=1,3
      NS=NIDS(I)
      110 READ (NREAD,825) (NID(I,J),J=1,NS)
      120 READ (NREAD,840) (VAF(I),I=1,NBODY)
      READ (NREAD,825) (NJNT(I),I=1,NPS)

      IF IOPT(3)=1, APPLY LINEARIZED BOUNCARY CCNDITION ON CHORDLINE
      OTHERWISE APPLY NONLINEAR BOUNDARY CONDITIONS ON AIRFOIL SURFACE

      IF (IOPT(3) .NE. 1) GO TO 116
      DO 115 J=1,NBODY
      I=NID(3,J)
      115 Y(I)=0.0
      116 CCNTINUE
      200 WRITE (NWRITE,920) NELS,NPS,NBW
      WRITE (NWRITE,930)
      CC 220 N=1,NELS
      220 WRITE (NWRITE,825) N,(NOD(N,J),J=1,4)
      WRITE (NWRITE,935)
      DO 230 I=1,NPS
      230 WRITE (NWRITE,940) I, NJNT(I),X(I),Y(I)
      WRITE (NWRITE,951) (NID(1,I),I=1,NFARF)
      WRITE (NWRITE,952) (NID(2,I),I=1,NWAKE)
      WRITE (NWRITE,953) (NID(3,I),I=1,NBODY)
      WRITE (NWRITE,955) (VAF(I),I=1,NBODY)

      REDEFINE MESH DATA, ECT. USING NEW NODAL NUMBERING SYSTEM
C
C
C
      CC 238 N=1,NELS
      DO 238 I=1,4
      IF (NOD(N,I) .EQ.0) GO TO 238
      KK=NOD(N,I)
      NCD(N,I)=NJNT(KK)
      CCNTINUE
      238 DO 239 I=1,3
      IS=NIDS(I)
      DO 239 J=1,IS
      KK=NID(I,J)
      239 NID(I,J)=NJNT(KK)

```

TRA000480
 TRA000490
 TRA000500
 TRA000510
 TRA000520
 TRA000530
 TRA000540
 TRA000550
 TRA000560
 TRA000570
 TRA000580
 TRA000590
 TRA000600
 TRA000610
 TRA000620
 TRA000630
 TRA000640
 TRA000650
 TRA000660
 TRA000670
 TRA000680
 TRA000690
 TRA000700
 TRA000710
 TRA000720
 TRA000730
 TRA000740
 TRA000750
 TRA000760
 TRA000770
 TRA000780
 TRA000790
 TRA000800
 TRA000810
 TRA000820
 TRA000830
 TRA000840
 TRA000850
 TRA000860
 TRA000870
 TRA000880
 TRA000890
 TRA000900
 TRA000910
 TRA000920
 TRA000930
 TRA000940
 TRA000950


```

C 243 I=1,NPS
RMLP(I)=X(I)
243 RML(I)=Y(I)
DC 244 II=1,NPS
I=NJNT(II)
X(I)=RMLP(II)
244 Y(I)=RML(II)
C
C HALF BANDWIDTH AND NUMBER OF EQUATIONS
C
NFBW=NBW/2
NEQ=3*NPS
C
C READ NONZERO INITIAL GUESS OR PROCEED WITH ZERO SOLUTION
C
IF (IOPT(2) .NE. 1) GO TO 250
READ (1,840) (S(I,NBW),I=1,NEQ)
GC TO 295
250 DC 260 I=1,NEQ
260 SLP(I)=0.0
C
C ITERATIONS START HERE AND CHECKED IF ITGIV IS EXCEEDED. IF SO,
C PRINT FAIL TO CONVERGE AND PROCEED TO NEXT CASE. OTHERWISE
C CONTINUE TO ITERATE
C
IRES=IRES+1
265 CCNTINUE
IF (IRES .GT. ITGIV) GO TO 600
FORMULATE SYSTEM OF ALGEBRAIC EQUATIONS
DC 266 I=1,NEQ
DC 266 J=1,NBW
266 S(I,J)=0.0
NMAT=1
CALL NEWK(SQMAC,NRM,NCM,NEQ,NBW,NEM,NELS,NOC,SLP,S,CCF,NPM,
1 X,Y,NMAT)
1 IMPCSE B.C. FOR FARFIELD, LINE OF SYMMETRY, AND ON AIRFCIL
IF (IOPT(4) .NE. 0) CALL BNDRY (S,UB,VB,VA,F,NID,NFARF,NBODY,
1 NWAKE,NBW,NHBW,NEQ,IOPT,IRES)
1 IF (IOPT(4) .NE. 0) GO TO 276
DC 274 I=1,NFARF
IE=3*NID(1,I)-3
DC 272 II=1,3
IE=IE+1
DC 270 K=1,NBW
S(IE,K)=0.0
270 S(IE,NHBW)=1.0
272 CCNTINUE
274 CCNTINUE
276 CONTINUE

```

TRA00960
 TRA00970
 TRA00980
 TRA00990
 TRA01000
 TRA01010
 TRA01020
 TRA01030
 TRA01040
 TRA01050
 TRA01060
 TRA01070
 TRA01080
 TRA01090
 TRA01100
 TRA01110
 TRA01120
 TRA01130
 TRA01140
 TRA01150
 TRA01160
 TRA01170
 TRA01180
 TRA01190
 TRA01200
 TRA01210
 TRA01220
 TRA01230
 TRA01240
 TRA01250
 TRA01260
 TRA01270
 TRA01280
 TRA01290
 TRA01300
 TRA01310
 TRA01320
 TRA01330
 TRA01340
 TRA01350
 TRA01360
 TRA01370
 TRA01380
 TRA01390
 TRA01400
 TRA01410
 TRA01420
 TRA01430


```

DC 280 I=1,NWAKE
IE=3*NID(2,I)
DC 278 K=1,NBW
S(IE,K)=0.0
278 S(IE,NHBW)=1.0
280 DC 285 J=1,NBODY
IE=3*NID(3,J)
DC 282 K=1,NBW
S(IE,K)=0.0
282 S(IE,NHBW)=1.0
S(IE,NHBW)=1.0
IF (IOPT(3) .NE. 1) S(IE,NHBW-1)=-VAF(J)
285 S(IE,NBW)=VAF(J)
CALL BANDED SOLVER TO SOLVE THE SYSTEM OF EQUATIONS
C
C
C
C
STORE STIFFNESS MATRIX AND INITIALIZE STARTING TIME DEPENDENT
UNKNOWN
CALL STCRE(S,WORK,NRM,NCM,NHCM,NFBW,NEQ,NSK)
IF (IRES .GT. 1 .OR. RMAC .GT. .91) GO TO 401
CALL BNCEQ(S,NRM,NCM,NEQ,NHBW)
DC 420 I=1,NRM
UT(I,3)=0.0
420 UT(I,2)=.5*S(I,NBW)
UT(I,1)=S(I,NBW)
GO TO 295
C
ASSEMBLE AND STORE THE DAMPING MATRIX
DC 401 DO 400 I=1,NRM
DC 400 J=1,NCM
400 S(I,J)=0.0
NMAT=2
CALL NEWK(SQMAC,NRM,NCM,NEQ,NBW,NEM,NELS,NOC,SLP,S,CCF,NPM,
1 X,Y,NMAT)
CALL STCRE(S,WORK,NRM,NCM,NHCM,NFBW,NEQ,NSC)
C
ASSEMBLE AND STORE THE MASS MATRIX
DO 410 I=1,NRM
DO 410 J=1,NCM
410 S(I,J)=0.0
NMAT=3
CALL NEWK(SQMAC,NRM,NCM,NEQ,NBW,NEM,NELS,NOC,SLP,S,CCF,NPM,
1 X,Y,NMAT)
CALL STCRE(S,WORK,NRM,NCM,NHCM,NFBW,NEQ,NSM)
C
CALL TIME(S,UT,WCRK,NRM,NCM,NHCM,NFBW,NEQ)
PRINT COMPUTED RESULTS
C
C
C
295 WRITE (NWRITE,970) RMAC,RFT
WRITE (NWRITE,975) IRES
CC 305 I=1,NPS
TRA01440
TRA01450
TRA01460
TRA01470
TRA01480
TRA01490
TRA01500
TRA01510
TRA01520
TRA01530
TRA01540
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J=NJNT(I)
II=3*NJNT(I)
PCT=S(II-2,NBW)
UPT=S(II-1,NBW)
V=S(II,NBW)
U=UPT+1.0
QSQ=U*U+V*V
ASQ=CONST*(1.0-QSQ)+1./SQMAC
RML(J)=SQRT(QSQ/ASQ)
PRATIO=(1.0+CONST*RML(J)**2)**EXP
CP=-2.*UPT
COF(J)=SQMAC*(1.0+2.4*UPT)
DELM=RML(J)-RMLP(J)
305 WRITE (NWRITE,978) I,POT,U,V,COF(J),RML(J),FRATIO,CP,DELM
C
C
C   DISPLAY PRESSURE COEFFICIENT CP
C
C   WRITE (NWRITE,985)
C   ISTOP=0
C   IFNT=0
C   DO 320 J=1,NBODY
C   IF (J.EQ. NBODY) ISTOP =1
C   I=NIID(3,J)
C   LPT=S(3*I-1,NBW)
C   V=S(3*I,NBW)
C   CF=-2.*UPT
C   CCNTINUE
320 IF (IRES .LT. 2) GO TO 382
C
C   CHECK CONVERGENCE. IF SO, PUNCH CCNVERGED SOLUTION AND
C   PROCEED TO NEXT CASE. OTHERWISE, UPDATE SOLUTION AND CONTINUE
C   TC ITERATE
C
C   DO 340 I=1,NPS
C   PCTE=1.0-RMLP(I)/RML(I)
C   IF (ABS(PCTE) .LT. ZTEST) GO TO 340
C   RFT=F1
C   GO TO 382
340 CCNTINUE
C   WRITE (NPUNCH,840) (S(I,NBW),I=1,NEQ)
C   GO TO 100
382 DO 385 I=1,NPS
385 RMLP(I)=RML(I)
C   RFTC=1.0-RFT
C   CC 390 I=1,NEQ
390 SLP(I)=RFT*S(I,NBW)+RFTC*SLP(I)
C   GO TO 265
600 WRITE (NWRITE,980)
  
```



```

805 WRITE (1,840) (S(I,NBW),I=1,NEQ)
820 GC TO 100
825 FCRMAT (18A4)
830 FCRMAT (40I2)
835 FCRMAT (16I5)
840 FCRMAT (15,3F10.0,15)
840 FCRMAT (1P8E10.3)
910 FCRMAT (1H1,2X,18A4//, CONVERGENCE LIMIT =',F6.4//)
920 FCRMAT (1H0,NO, OF ELEMENTS=',I4,', NO. CF NODES=',I4,
1 FCRMAT (, FULL BANDWIDTH =',I4//)
930 FCRMAT (OELE,NO, AND ELEMENT NODES')
935 FCRMAT (OOLD NODE,, NEW NODE',6X,X(I)',Y(I)')
940 FCRMAT (16,110,2E15.5)
951 FCRMAT (ONODES AT FAR FIELD',(20I5))
952 FCRMAT (ONODES ON THE LINE OF SYMMETRY',(20I5))
953 FCRMAT (ONODES ON THE AIRFOIL',(20I5))
955 FCRMAT (OSLOPE ALONG NODES ON AIRFOIL',(8E15.5))
970 FCRMAT (1H1,MACH NUMBER=',F6.3,5X,RELAX. FACTOR =',F6.4//)
975 FCRMAT (1H0,4X,NO. OF ITERATIONS =',I4/
1 FCRMAT (1H0,7X,NO. OF ITERATIONS =',I4/
2 FCRMAT (1H0,7X,NO. OF ITERATIONS =',I4/
980 FCRMAT (12X,COF,11X,LMAC,8X,PHIT,11X,UCOM,11X,VCCM,
978 FCRMAT (12X,COF,11X,LMAC,8X,PHIT,11X,UCOM,11X,VCCM,
985 FCRMAT (12X,COF,11X,LMAC,8X,PHIT,11X,UCOM,11X,VCCM,
2000 FCRMAT (10,1P8E15.4)
END

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TRA02700

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SUBROUTINE BCOND

```

SUBROUTINE BCOND(U,V,IOPT,N)
DIMENSION U(N),V(N),UP(50),VP(50)
DATA NWRITE/6/
DO 10 I=1,N
  UP(I)=1+U(I)
  VP(I)=V(I)
CCNTINUE
10 IF (IOPT.EQ.1) GO TO 150
  WRITE(NWRITE,100)
100 FCRMAT( , UCOM AT INLET.)
  N2=N/2
  N3=N2+1
  WRITE (NWRITE,110) (UP(I),I=1,N2)
  FCRMAT ( , VCOM AT INLET.)
120 WRITE (NWRITE,110) (VP(I),I=1,N2)
  WRITE (NWRITE,130)
  FCRMAT ( , UCOM AT EXIT.)
130 WRITE (NWRITE,110) (UP(I),I=N3,N)
  WRITE (NWRITE,140)
  FCRMAT ( , VCOM AT EXIT.)
140 FCRMAT ( , UCOM AT EXIT.)
110 FCRMAT ( , VCOM AT EXIT.)
  WRITE (NWRITE,110) (VP(I),I=N3,N)
  GO TO 160
150 WRITE(NWRITE,121)
  WRITE(NWRITE,110) (UP(I),I=1,N)
  WRITE(NWRITE,122)
  WRITE(NWRITE,110) (VP(I),I=1,N)
  FCRMAT( , UCOM AT SONIC LINE.)
121 FCRMAT( , VCOM AT SONIC LINE.)
122 FCRMAT( , VCOM AT SONIC LINE.)
160 RETURN
END

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SUBROUTINE BNDEQ

```

1201 SUBROUTINE BNDEQ(A,NRMAX,NCMAX,N,ITERM)
C      CCNTINUE
C      EQUATION SOLVER FOR BANDED NON-SYMMETRIC SYSTEM OF EQUATIONS
C      SOLUTION STORED IN THE LAST COLUMN AT (I,2*ITERM)
C
C      DIMENSION A(NRMAX,NCMAX)
C      DATA NREAD,NWRITE,NPUNCH/4,8,3/
C      CERO=1.E-6
C      PARE=CERO**2
C      NBND=2*ITERM
C      NEM=NBND-1
C
C      BEGINS ELIMINATION OF THE LOWER LEFT
C
C      DO 1000 I=1,N
C      IF (ABS(A(I,ITERM))) .LT. CERO) GO TO 410
C      GO TO 430
C      IF (ABS(A(I,ITERM))) .LT. PARE) GO TO 1600
C      410 WRITE (6,420) A(I,ITERM), I
C      420 FORMAT (10,420) A(I,ITERM), I
C      430 JLAST=MINO(I+ITERM-1,N)
C      L=ITERM+1
C      CCNTINUE
C      DO 500 J=I,JLAST
C      L=L-1
C      IF (ABS(A(J,L))) .LT. PARE) GO TO 500
C      B=A(J,L)
C      DO 450 K=L,NBND
C      A(J,K)=A(J,K)/B
C      450 IF (I .EQ. N) GO TO 1200
C      CCNTINUE
C      L=0
C      JFIRST=I+1
C      IF (JLAST .LE. I) GO TO 1000
C      DO 900 J= JFIRST,JLAST
C      L=L+1
C      IF ( ABS(A(J,ITERM-L)) .LT. PARE) GO TO 900
C      DO 600 K=ITERM,NBM
C      A(J,K-L) = A(J-L,K) -A(J,K-L)
C      600 A(J,NBND) = A(J-L,NBND) - A(J,NBND)
C      IF (I .GE. N-ITERM+1) GO TO 900
C      DO 800 K=1,L
C      A(J,NBND-K) = -A(J,NBND-K)
C      800 CCNTINUE
C      900 CONTINUE
C      1000

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C      BACK-SUBSTITUTION
C      120C  L=ITERM -1
C      DO 1500 I=2,N
C      DO 1500 J=1,L
C      IF (N+1-I+J.GT. N) GO TO 1500
C      ATEMP1=A(N+1-I,NBND)
C      ATEMP2=A(N+1-I+J,NBND)
C      ATEMP3=A(N+1-I,ITERM+J)
C      A(N+1-I,NBND)=A(N+1-I,NBND)-A(N+1-I+J,NBND)*A(N+1-I,ITERM+J)
C      1500 CCCONTINUE
C      RETURN
C      PRINT THE ENTIRE MATRIX IF ZERO CN MAIN DIAGCNAL
C      1600 WRITE (6,1601)
C      1601 FORMAT (10X,COMPUTATION STOPED IN BNDEQ BECAUSE ZERC APPEARED ON
C      1602 1,MAIN DIAGONAL. THE MATRIX FOLLOWS.)
C      DO 1602 I=1,N
C      1602 WRITE (NWRITE,1603) (A(I,J), J=1,NBND)
C      1603 FCRMAT (10E12.4)
C      STOP
C      END

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TRA00690
TRA00700

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SUBROUTINE BNDRY

```

SUBROUTINE BNDRY(S,UB,VB,VAF,NID,NFARF,NBODY,NWAKE,NBW
1,NFBW,NEQ,IOPT,IRES)
DIMENSION S(456,60),UB(50),VB(50),VAF(50),NID(3,50),IOPT(40)
NBW$=NBW-1
IF (IOPT(5) .EQ. 1 .AND. IRES .GT. 1) NFARF=NFARF/2
DC 10 I=1,NFARF
IE=3*NID(1,I)-2
DC 10 J=1,2
IE=IE+1
NX=IE-NHBW
BC=UB(I)
IF (J.EQ.2) BC=VB(I)
DO 41 L=1,NBW$
LX=NX+L
IF (LX .LE. 0 .OR. LX .GT. NEQ) GO TO 41
S(LX,NBW)=S(LX,NBW)-BC*S(LX,NBW-L)
S(LX,(NBW-L))=0.
CCNTINUE
DC 30 K=1,NBW
S(IE,K)=0.0
CCNTINUE
S(IE,NHBW)=1.0
S(IE,NBW)=BC
DC 100 I=1,NBODY
IE=3*NID(3,I)
BC=VAF(I)
NX=IE-NHBW
DC 410 L=1,NBW$
LX=NX+L
IF (LX .LE. 0 .OR. LX .GT. NEQ) GO TO 410
S(LX,NBW)=S(LX,NBW)-BC*S(LX,NBW-L)
S(LX,(NBW-L))=0.0
CCNTINUE
DC 300 K=1,NBW
S(IE,K)=0.0
S(IE,NHBW)=1.0
S(IE,NBW)=BC
DC 200 I=1,NWAKE
IE=3*NID(2,I)
DC 330 K=1,NBW
S(IE,K)=0.0
S(IE,NHBW)=1.
IF (IOPT(5) .EQ. 1 .AND. IRES .GT. 1)NFARF=NFARF*2
RETURN
END

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SUBROUTINE EMTC

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SUBROUTINE EMTC(A,AT,ATT,XL,YL,PEL,SQMAC)
EVALUATE ELEMENT MATRIX FOR A TRIANGLE BY GAUSSIAN QUADRATURE
SUBROUTINE DERV CALLED TO EVALUATE SHAPE FUNCTION DERIVATIVES
AT THE GAUSSIAN POINTS

DIMENSION A(9,9),P(9),Q(9),NP(5),B(3),C(3),XL(3),YL(3),S(3),
1 DNX(9),DNXX(9),DNY(9),PEL(9),
2 DN(9),AT(9,9),ATT(9,9)
DIMENSION EINT(3,7),WT(7)
DATA LMAX/7/,WT/0.225,3*0.13239415,3*.12593518/
DATA EINT/3*0.33333333,0.05961587,3*0.47C14206,0.05961587,
1 3*0.47C14206,0.05961587,0.79742699,3*0.1C128651,
2 0.79742699,3*0.10128651,0.79742699/
DATA NP/1,2,3,1,2/,GAMMA/1.40/
DATA NREAD,NWRITE,NPUNCH/4,8,3/
DO 2 I=1,9
CC 2 J=1,9
ATT(I,J)=0.
2 A(I,J)=0.0
CC 4 I=1,3
J=NP(I+1)
K=NP(I+2)
B(I)=YL(J)-YL(K)
C(I)=XL(K)-XL(J)
AREA=0.5*(B(2)*C(3)-B(3)*C(2))
CST1=1.0-SQMAC
CST2=SQMAC*(1.0+GAMMA)
DO 10 L=1,LMAX
DO 10 I=1,3
DO 10 J=1,3
S(I)=EINT(I,L)
10 CALL DERV (AREA,B,C,S,DN,DNX,DNXX,DNYY)
U=0.0
UX=0.0
DO 30 I=1,9
U=U+DNX(I)*PEL(I)
30 UX=UX+DNXX(I)*PEL(I)
ALPHA=CST1-CST2*U
DO 40 I=1,9
P(I)=ALPHA*DNXX(I)+DNY(I)
40 Q(I)=P(I)-CST2*UX*DNX(I)
WEIGHT=WT(L)*AREA
CC 60 I=1,9
CST=WEIGHT*Q(I)

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 TRA00450


```

DO 60 J=1,9
  AT(I,J)=AT(I,J)-2*CST*SQMAC*DNX(J)
  ATT(I,J)=ATT(I,J)-CST*SQMAC*DN(J)
  A(I,J)=A(I,J)+CST*P(J)
60 CCNTINUE
1CC RETURN
END

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TRA00460
TRA00470
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TRA00510
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SUBROUTINE EMQT

```

C
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C
SUBROUTINE EMQT(XQ,YQ,PMQ,SQMAC,EG,EQT,ETT,NTRS)
    GENERATE MATRIX FOR A QUADRILATERAL OR TRIANGLE
    SUBROUTINE EMTC CALLED TO GENERATE MATRIX FOR A BASIC TRIANGLE

    DIMENSION EQ(12,12),ET(9,9),XQ(4),YQ(4),XT(3),YT(3),MP(3,4)
    DIMENSION PMQ(12),PMT(9),EQI(12,12),EQTI(12,12),ETI(9,9),
1    ETIT(9,9)
    DATA MP/1,2,3,3,4,1,2,3,4,4,1,2/
    DATA NREAD,NWRITE,NPUNCH/4,8,3/
    FTOR=1.0
    IF (NTRS.EQ. 4) FTOR=.5
    DO 100 I=1,12
    CC 100 J=1,12
    EQT(I,J)=0.0
    EQTI(I,J)=0.0
    EQ(I,J)=0.0
10C
    DO 150 II=1,NTRS
    CC 105 I=1,3
    NI=MP(I,II)
    IT=3*(I-1)
    IQ=3*(NI-1)
    DO 102 J=1,3
    IT=IT+1
    IC=IQ+1
102
    PMT(IT)=PMQ(IQ)
    XT(I)=XQ(NI)
    YT(I)=YQ(NI)
105
    CALL EMTC(ET,ETT,ETTT,XT,YT,PMT,SQMAC)
    DO 130 K=1,3
    NR=3*(MP(K,II)-1)
    IE=3*(K-1)
    DO 130 KK=1,3
    NR=NR+1
    IE=IE+1
    CC 130 L=1,3
    NC=3*(MP(L,II)-1)
    JE=3*(L-1)
    DO 130 LL=1,3
    NC=NC+1
    JE=JE+1
    EQT(NR,NC)=EQT(NR,NC)+ETT(IE,JE)*FTOR
    EQTI(NR,NC)=EQTI(NR,NC)+ETTT(IE,JE)*FTOR
120
    EC(NR,NC)=EQ(NR,NC)+ET(IE,JE)*FTOR
150
    CC CONTINUE

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 TRA00450

RETURN
END

TRAC0460
TRA00480

SUBROUTINE DERV

```

C
C
C
SUBROUTINE DERV(AREA,B,C,S,DN,DNX,DNXX,DNYY)
EVALUATE SHAPE FUNCTION DERIVATIVES AT GAUSSIAN PCINT
DIMENSION B(3),C(3),S(3),DN(9),DNX(9),DNXX(5),DNYY(9),NP(5)
DATA NP/1,2,3,1,2/
CATA NREAD,NWRITE,NPUNCH/4,8,3/
TWCA=2.*AREA
TWCASQ=TWCA**2
CC 200 I=1,3
J=NP(I+1)
K=NP(I+2)
SI=S(I)
SJ=S(J)
SK=S(K)
BI=B(I)
BJ=B(J)
BK=B(K)
CI=C(I)
CJ=C(J)
CK=C(K)
SISQ=SI*SI
BISQ=BI*BI
CISQ=CI*CI
ALFA=0.5*(CK-CJ)
BETA=0.5*(BJ-BK)
F=SI*SJ*SK
FX=BI*SJ*SK+BJ*SK*SI*SJ
FXX=2.*(SI*BJ*SK+SJ*BK*BI+SK*BI*BJ)
HYY=2.*(SI*CJ*CK+SJ*CK*CI+SK*CI*CJ)
CSS=6.*(SI*(1.0-SI)
CS=6.*(1.0-2.0*SI)
LN=3*I-2
DN(L)=SI*(3-2*SI)+2*H
DNX(L)=BI*CSS+2.*HX
DNXX(L)=BISQ*CS+2.*HXX
DNYY(L)=CISQ*CS+2.*HYY
CS=CK*SJ-CJ*SK
L=L+1
DN(L)=SISQ*CS+ALPHA*H
DNX(L)=2.*BI*SI*CS+TWCA*SI*SQ+ALFA*HX
DNXX(L)=2.*BI*SI*CS+4.*BI*TWCA*SI+ALFA*HXX
DNYY(L)=2.*CISQ*CS+ALFA*HYY
ES=BJ*SK-BK*SJ
L=L+1
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DN(L)=SISQ*BS+BETA*H
CNX(L)=2.*BI*SI*BS+BETA*HX
CNXX(L)=2.*BISQ*BS+BETA*HXX
CNY(L)=2.*CISQ*BS+4.*CI*TWOA*SI+BETA*HYY
200 CC 300 I=1,9
    CNX(I)=DNX(I)/TWOA
    CNXX(I)=DNXX(I)/TWOASQ
300 DNY(I)=DNY(I)/TWOASQ
    RETURN
    ENC

```


SUBROUTINE NEWK

```

C
C
C
C
C
SUBROUTINE NEWK(SQMAC,NRM,NCM,NEQ,NBW,NEM,NELS,NOD,SLP,S,
1 COEF,NPM,X,Y,NMAT)
GENERATE SYSTEM MATRIX BY ASSEMBLING CONTRIBUTIONS FROM
ALL THE ELEMENTS
SUBROUTINE EMQT CALLED TO GENERATE ELEMENT MATRIX
C
DIMENSION COEF(NPM),X(NPM),Y(NPM),XQ(4),YQ(4),PM(12),BB(12,12)
DIMENSION BBC(12,12),BBM(12,12)
DIMENSION NOD(NEM,4),S(NRM,NCM),SLP(NRM)
DATA NREAD,NWRITE,NPUNCH/4,8,3/
NFEW=NBW/2
CC 480 N=1,NELS
I1=1
IF (NOD(N,4)) 402,402,404
402 NPEL=3
NTRS=1
GC TO 410
404 NPEL=4
NTRS=4
I2=0
CC 408 I=1,4
NI=NOD(N,I) .GT. 1.00) I2=I2+1
408 IF (COEF(NI) .GT. 1.00) NTRS=2
IF (I2.EQ.0) NTRS=3
IF (I2.EQ.4) I1=3
41C DC 425 I=1,NPEL
NI=NOD(N,I)
XQ(I)=X(NI)
YQ(I)=Y(NI)
DC 425 J=1,3
IS=3*(NI-1)+J
IE=3*(I-1)+J
PM(IE)=SLP(IS)
425 CALL EMQT(XQ,YQ,PM,SQMAC,BB,BBC,BEM,NTRS)
DC 450 I=1,NPEL
NF=3*(NOD(N,I)-1)
IE=3*(I-1)
DC 450 II=1,3
NR=NR+1
IE=IE+1
DC 450 J=1,NPEL
NC=3*(NOD(N,J)-1)-NR+NHBW
JE=3*(J-1)
CC 450 JJ=1,3
TRA000010
TRA000020
TRA000030
TRA000040
TRA000050
TRA000060
TRA000070
TRA000080
TRA000090
TRA000100
TRA000110
TRA000120
TRA000130
TRA000140
TRA000150
TRA000160
TRA000170
TRA000180
TRA000190
TRA000200
TRA000210
TRA000220
TRA000230
TRA000240
TRA000250
TRA000260
TRA000270
TRA000280
TRA000290
TRA000300
TRA000310
TRA000320
TRA000330
TRA000340
TRA000350
TRA000360
TRA000370
TRA000380
TRA000390
TRA000400
TRA000410
TRA000420
TRA000430
TRA000440
TRA000450

```



```

NC=NC+1
JE=JE+1
IF (NMA T-2) 443,442,441
441 S(NR,NC)=S(NR,NC)+BBM(IE,JE)
GO TO 450
442 S(NR,NC)=S(NR,NC)+BBC(IE,JE)
GC TO 450
443 S(NR,NC)=S(NR,NC)+BB(IE,JE)
45C CCNTINUE
48C RETURN
ENC

```

```

TRA00460
TRA00470
TRA00480
TRA00490
TRA00500
TRA00510
TRA00520
TRA00530
TRA00540
TRA00550
TRA00560
TRA00580

```


SUBROUTINE STORE

```

SUBROUTINE STORE (S,WCRK,NRM,NCM,NF-CM,NHBW,NEQ,NDATA)
DIMENSION S(NRM,NCM),WORK(NRM,NHCM)
REWIND NDATA
1C  DO 10 I=1,NRM
    DO 10 J=1,NHCM
        WCRK(I,J)=0.0
        CC 20 I=1,NEQ
        DO 20 J=1,NHBW
            WCRK(I,J)=S(I,J)
            WRITE (NDATA) WORK
            CC 30 I=1,NEQ
            DO 30 J=1,NHBW
                JA=NHBW+J
                WCRK(I,J)=S(I,JA)
            WRITE (NDATA) WORK
        REWIND NDATA
    RETURN
END

```

STG000010
 STG000020
 STG000030
 STG000040
 STC000050
 STC000060
 STG000070
 STG000080
 STG000090
 STG000100
 STG000110
 STC000120
 STG000130
 STG000140
 STC000150
 STG000160
 STC000170
 STG000180

SUBROUTINE TIME

```

C
C
C
SUBROUTINE TIME(S,U,WORK,NRM,NCM,NFCM,NHBW,NEG)
HCLBOLT INTEGRATION FOR UNSTEADY PROBLEMS
DIMENSION S(NRM,NCM),WORK(NRM,NHCM),U(NRM,3),NDATA(3)
DATA NDATA/9,13,14/
CATA DELT/100./
NBW=2*NHBW
NBW1=NBW-1
AC=2./DELT**2
A1=11./(6*DELT)
A2=5./DELT**2
A3=3./DELT
A4=-2.*A0
A5=-A3/2.
A6=A0/2.
A7=A3/9
CC 100 I=1,NRM
CC 100 J=1,NCM
DO 100 S(I,J)=0.0
C FCRM LOWER TRIANGULAR PLUS DIAGONAL EFFECTIVE MATRIX
DO 25 K=1,3
NCATA$=NDATA(K)
REAC (NCATA$) WORK
IF (K-2) 1,2,3
1 CCNST=1
2 CCNST=A0
3 CCNST=A1
4 CCNST=A1
10 S(I,J)=WORK(I,J)*CONST+S(I,J)
C
C
C
MULTIPLY CONTRIBUTING MATRICES BY THE TIME CEPENDENT
VECTORS TO FORM RIGHT HAND SIDE
DO 20 J=1,NHBW
JM=NHBW+1-J
IM=0
CC 20 I=JM,NEG
IM=IM+1
IF (K-2) 11,12,13
11 UW=0.
12 UW=A2*U(IM,1)+A4*U(IM,2)+A6*U(IM,3)
CC TO 20
CC TO 20

```

TIM00010
 TIM00020
 TIM00030
 TIM00040
 TIM00050
 TIM00060
 TIM00070
 TIM00080
 TIM00090
 TIM00100
 TIM00110
 TIM00120
 TIM00130
 TIM00140
 TIM00150
 TIM00160
 TIM00170
 TIM00180
 TIM00190
 TIM00200
 TIM00210
 TIM00220
 TIM00230
 TIM00240
 TIM00250
 TIM00260
 TIM00270
 TIM00280
 TIM00290
 TIM00300
 TIM00310
 TIM00320
 TIM00330
 TIM00340
 TIM00350
 TIM00360
 TIM00370
 TIM00380
 TIM00390
 TIM00400
 TIM00410
 TIM00420
 TIM00430
 TIM00440
 TIM00450


```

C
12 UM=A3*U(IM,1)+A5*U(IM,2)+A7*U(IM,3)
20 S(IM,NBW)=S(IM,NBW)+WORK(I,J)*UM
25 CCNT=INUE
FCRM UPPER TRIANGULAR EFFECTIVE MATRIX
CC 45 K=1,3
NCATA$=NDATA(K)
READ (NCATA$) WORK
IF(K-2) 21,22,23
21 CCNST=1
CC TO 24
22 CCNST=A0
CC TO 24
23 CCNST=A1
CC 30 J=1,NHBW
24 DO 30 I=1,NEQ
JA=J+NHBW
S(I,JA)=S(I,JA)+WORK(I,J)*CONST
3C MAXJ=NHBW-1
CC 40 J=1,MAXJ
JA=NHBW+J
MAXI=NEQ-J
IM=J
DC 40 I=1,MAXI
IM=IM+1
IF (K-2) 31,32,33
31 UM=0
GC TO 40
32 UM=A2*U(IM,1)+A4*U(IM,2)+A6*U(IM,3)
GC TO 40
33 UM=A3*U(IM,1)+A5*U(IM,2)+A7*U(IM,3)
40 S(IM,NBW)=S(IM,NBW)+WORK(I,J)*UM
45 CCNT=INUE
843 FCRMAT(IH,4(F10.5,5X))
CALL BNDEQ(S,NRM,NCM,NEQ,NHBW)
CC 50 I=1,NRM
U(I,3)=U(I,2)
U(I,2)=U(I,1)
50 U(I,1)=S(I,NBW)
8CC FORMAT(IH,6(F10.5,2X))
RETURN
END

```

```

TIM00460
TIM00470
TIMC0480
TIM00490
TIM00500
TIM00510
TIM00520
TIM00530
TIM00540
TIM00550
TIM00560
TIM00570
TIM00580
TIM00590
TIM00600
TIM00610
TIM00620
TIM00630
TIM00640
TIM00650
TIM00660
TIM00670
TIM00680
TIM00690
TIM00700
TIM00710
TIM00720
TIM00730
TIM00740
TIM00750
TIM00760
TIM00770
TIM00780
TIM00790
TIM00800
TIM00810
TIM00820
TIM00830
TIM00840
TIM00850
TIM00870

```


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tion by the finite ele-
ment method.

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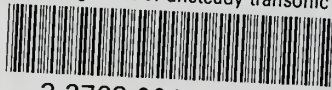
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